## ECE 145A / 218 C, notes set 2: Transmission Line Parasitics

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#### **Transmission Lines**

Approximate properties of microstrip line.

Skin Effect Losses

substrate modes and loss by coupling into these.

Lateral modes on lines

Excitation of unwanted circuit-like modes, ground continuity

Packaging and power supply resonances

## **Skin Loss**

### Skin effect losses l

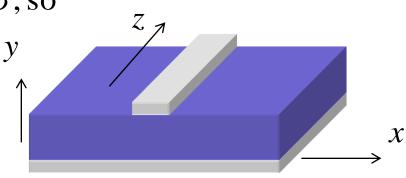
Wave equation inside metal (ignore *x* variation)  $k_y^2 + k_z^2 = k^2$  where  $k^2 = (j\omega\mu)(j\omega\varepsilon + \sigma)$ 

The wavelength along the transmission - line is long, and the penetration distance of current into the metal is small,

so 
$$k_y^2 >> k_z^2 \rightarrow k_y^2 = k^2 = (j\omega\mu)(j\omega\varepsilon + \sigma)$$

In a metal, at low frequencies  $j\omega\varepsilon \ll \sigma$ , so

 $k_y^2 = j\omega\mu\sigma$ 



### Skin effect losses II

In a metal, at low frequencies  $j\omega\varepsilon \ll \sigma$ , so

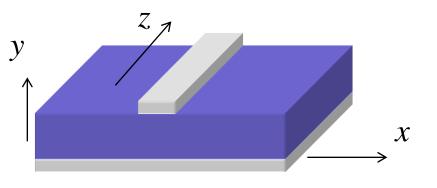
$$k_{y} = \pm \sqrt{j\omega\mu\sigma} = \pm (1+j)\left(\frac{\omega\mu\sigma}{2}\right)$$

Defining the \*skin depth\* as  $\delta = \sqrt{2/\omega\mu\sigma}$ :  $E(z) = E_o e^{-y/\delta} e^{-jy/\delta}$ 

The field dies down exponentially with distance into the metal. The (1/e) penetration depth is the skin depth  $\delta$ 

 $\delta$  varies as  $\omega^{-1/2}$ 

At 100GHz in Gold,  $\delta \cong 200 \text{ nm}$ 



### **Skin effect losses III**

Let us treat this approximately :

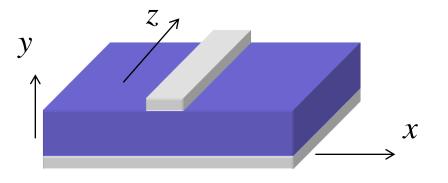
The conductor only carries current in a layer of thickness  $\delta$ .

With conductivity  $\sigma$  and width W,

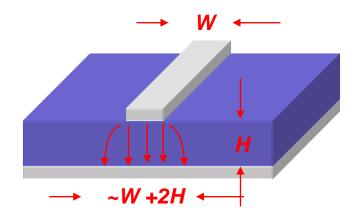
the conductor has resistance per unit length

 $R_{series} / L = 1 / \sigma \delta W = (1 / \sigma W) \cdot \sqrt{\omega \mu \sigma / 2} = (1 / W) \cdot \sqrt{\omega \mu / 2\sigma}$ 

A more careful treatment develops the concept of surface impedance. See the appendix.



#### Skin effect losses IV



Transmission - lines have skin effect in both the signal and ground conductors. For microstrip :

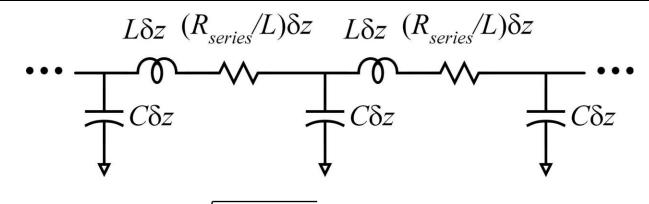
$$R_{series} / L \approx \frac{1}{\delta\sigma} \frac{1}{W} + \frac{1}{\delta\sigma} \frac{1}{W + 2H} = \left(\frac{1}{W} + \frac{1}{W + 2H}\right) \frac{1}{\delta\sigma} = \frac{1}{P\delta\sigma}$$

In general, we can write this as

$$R_{series} / L = 1 / \delta \sigma P = (1 / P) \cdot \sqrt{\omega \mu / 2\sigma}$$

where *P* is the effective current - carrying periphery.

#### **Skin effect losses IV**



$$R_{series} / L = (1 / P) \cdot \sqrt{\omega \mu / 2\sigma}$$

From our earlier transmission - line analysis, this introduces attenuation per unit distance

$$\alpha \cong \frac{R_{series} / L}{2Z_0} \quad \alpha \propto \sqrt{\omega}$$

This is called skin loss

## Loss Tangent

#### **Loss Tangent**

#### Common dielectrics also introduce high - frequency attenuation.

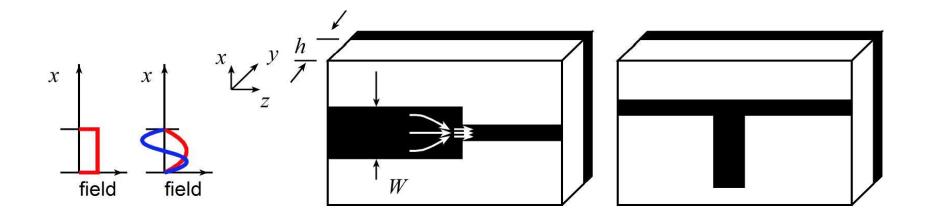
This effect is quantified by a \*loss tangent \*  $\varepsilon_{\rm r} = \varepsilon_{\rm r,real} + j\varepsilon_{\rm r,imaginary} = \varepsilon_{\rm r,real} (1 + j \tan(\delta))$ 

We should be aware of dielectric losses,

but we will not discuss these further in this class.

# transverse transmission-line modes

#### **Lateral Modes (1)**



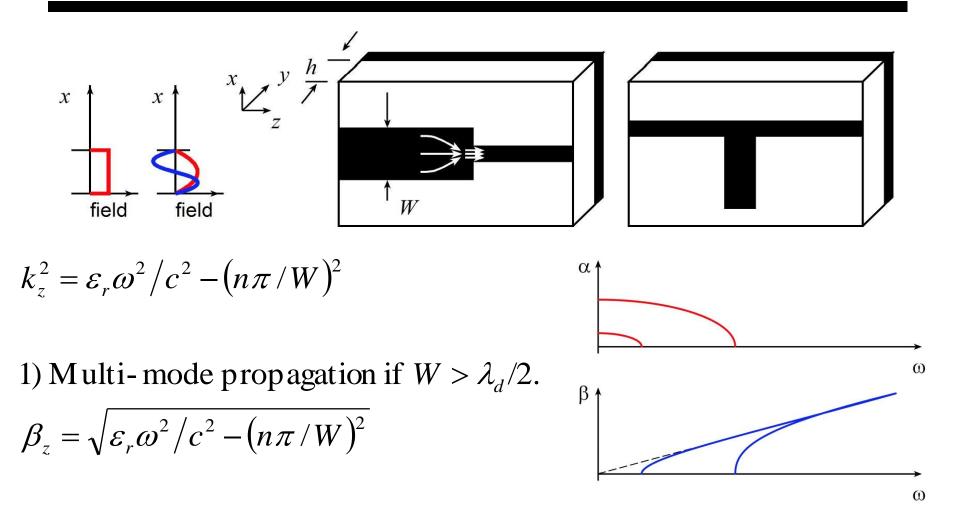
•

In dielectric : waves of form 
$$\vec{E}_0 e^{j\omega t} e^{\pm jk_x x} e^{\pm jk_y y} e^{\pm jk_z z}$$
  
 $k_x^2 + k_y^2 + k_z^2 = k^2 = \varepsilon_r \omega^2 / c^2 = (2\pi/\lambda_d)^2$ 

Waves can propagate\*laterally \* on transmission - line :

$$k_y = 0$$
 and  $k_x = n\pi/W$  for  $n = 0, 1, 2, ...$   
 $\rightarrow k_z^2 = \varepsilon_r \omega^2 / c^2 - (n\pi/W)^2$ 

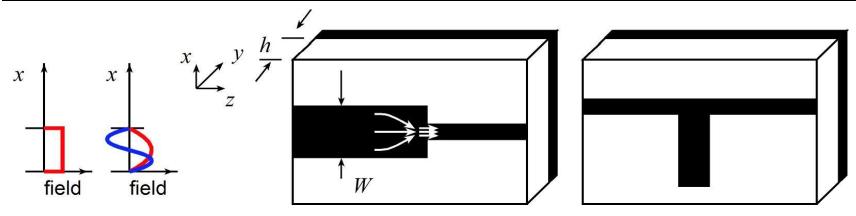
#### Lateral Modes (2)



2) Evanescent propagation  $e^{-\alpha_z z}$  if  $W < \lambda_d/2$ :

$$\alpha_z = \sqrt{(n\pi/W)^2 - \varepsilon_r \omega^2/c^2}$$

#### **Lateral Modes---and Junction Parasitics (3)**



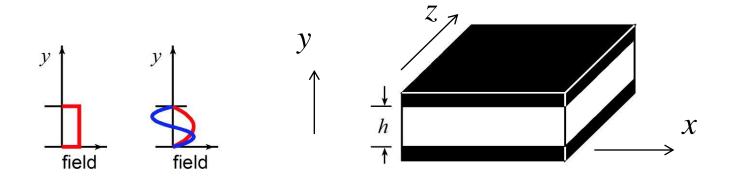
Evanescent propagation  $e^{-\alpha_z z}$  if  $W \cong \lambda_d/2$ :

Reactive power in evanescent modes → junction parasitics ADSlibrary junction models, or electromagnetic simulation. Lessons:

lines must be much narrower than a half - wavelength. must model junction parasitics

# Substrate Modes and Radiation Loss

#### **Substrate Modes**



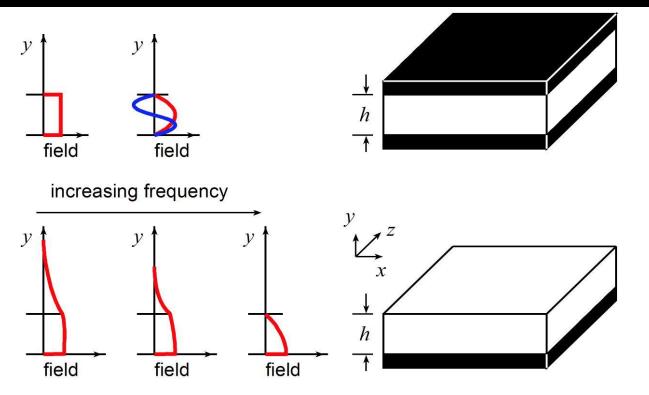
In dielectric : waves of form  $E_0 e^{j\omega t} e^{\pm jk_x x} e^{\pm jk_y y} e^{\pm jk_z z}$ 

$$k_x^2 + k_y^2 + k_z^2 = k^2 = \varepsilon_r \omega^2 / c^2 = (2\pi/\lambda_d)^2$$

We can have standing waves across the substrate thickness :

$$k_y = n\pi / h$$
 for  $n = 0, 1, 2, ...$ 

#### **Substrate Modes**



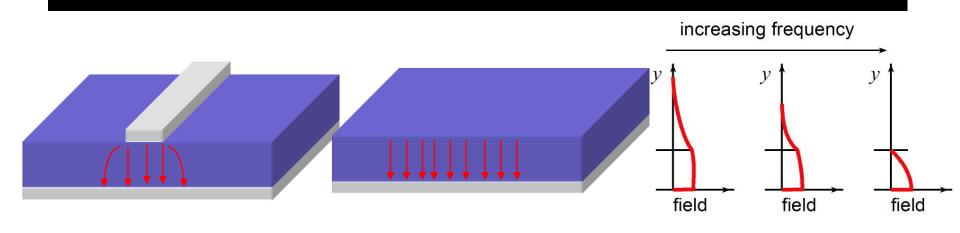
Substrate with top, bottom metal surfaces

$$\rightarrow$$
 modes with  $h = \lambda_d / 2, \lambda_d, 3\lambda_d / 2...$ 

Substrate with no top metal  $\rightarrow$  tranverse *E* - mode;

strongly confined as  $\lambda_d/4 \rightarrow T$ ; weakly confined at low frequencies.

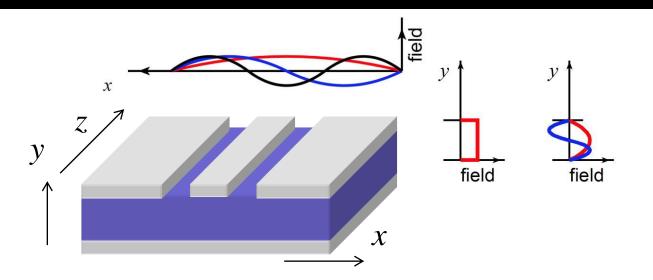
#### **Substrate Mode Coupling: Microstrip**



These dielectric slab modes can propage in x and in z.

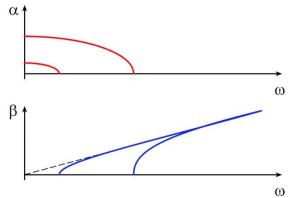
Nonzero mode coupling ("radiation") loss at all frequencies. Very strong mode coupling when  $h \ge \lambda_d/4$ 

#### **Substrate Mode Coupling: CPW**

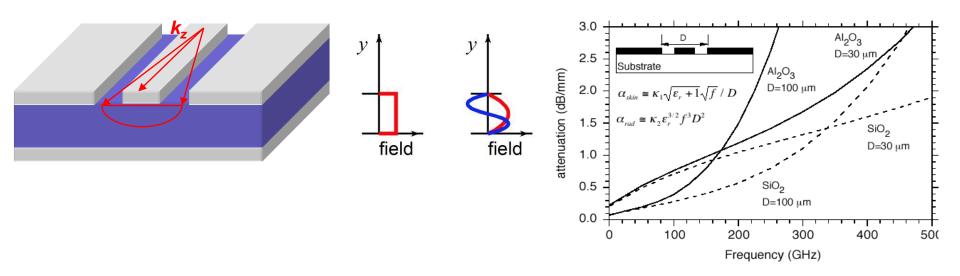


Substrate modes are allowed when  $\lambda_d \leq 2 \cdot (\text{substrate} \text{thickness})$ Waves are of form  $\vec{E}_0 e^{j\omega t} e^{\pm jk_x x} e^{\pm jk_y y} e^{\pm jk_z z}$ 

 $\rightarrow k_x^2 + k_y^2 + k_z^2 = k^2$  where  $k^2 = c^2 / \varepsilon_r \omega^2 = (2\pi/\lambda_d)^2$  and  $k_y = (n\pi/h)$ 



#### **Substrate Mode Coupling: CPW**

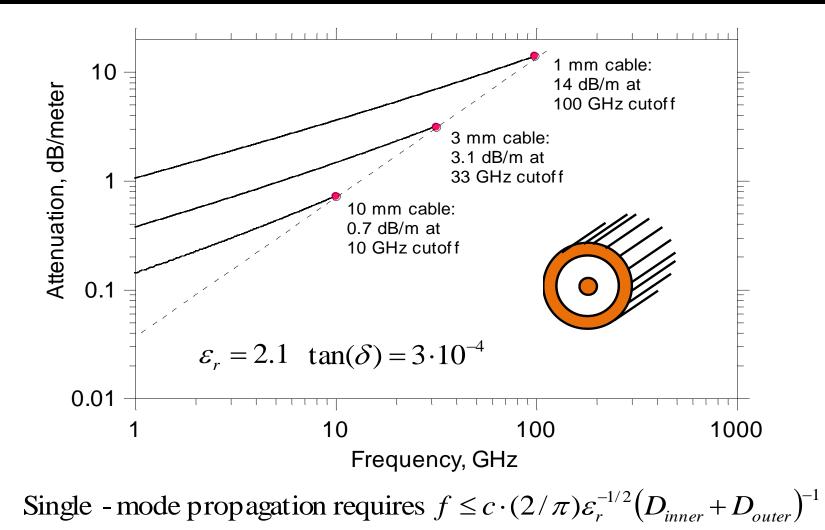


- Modes couple strongly when  $k_{y,CPW} = k_{y,substratemode}$ Given thick substrate,  $H >> \lambda_d$ :
- mode coupling loss, dB/mm  $\propto$  (line transverse dimensions)<sup>2</sup> · frequency <sup>2</sup> "radiation loss"

If we use narrow lines and thin substrates then skin - effect losses will be large.

If we use wide lines and thick substrates then lateral modes and substrateradiation will be major problems.

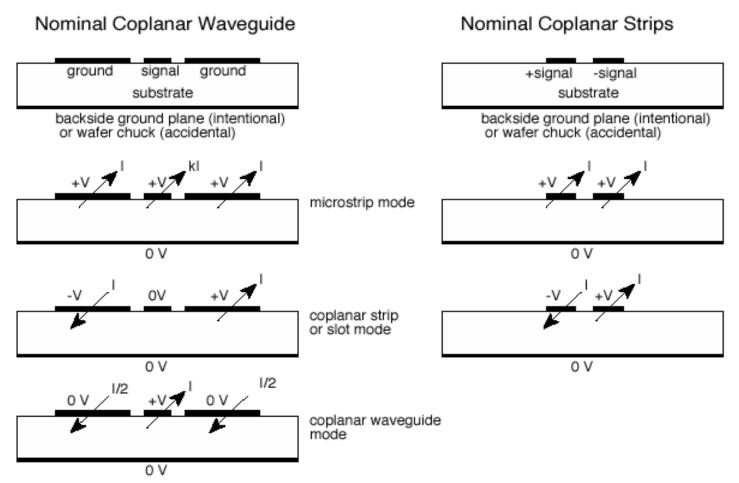
#### **Loss of Coaxial Cable**



Skin loss  $\alpha_{skin} \propto f^{1/2} / D_{inner} \longrightarrow \text{Loss } \alpha_{skin} \propto f^{3/2}$ 

# "circuit-type" parasitic modes

### **Transmission-Line Parasitic Modes**



- Total number of quasi-TEM modes is one less than # of conductors
- · Care must be taken to avoid excitation of parasitic modes
- unexpected results will otherwise arise...

#### To Avoid "Circuit-Type" Parasitic Modes

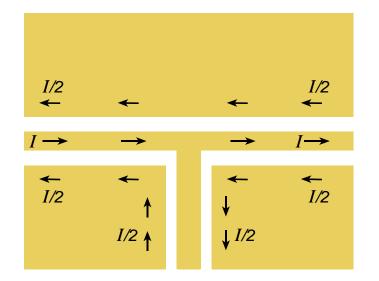
- 1) Where do the currents flow?
- 2) Which conductors have what voltages for which modes?

Be aware that:

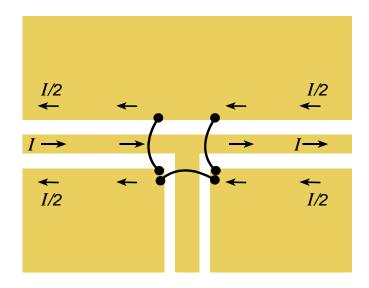
 currents must flow in the ground planes of unbalanced transmission lines. The currents flow close to the edge of the ground plane nearest the signal conductor.

 there are equal and opposite voltages on the 2 conductors of balanced transmission lines. This seriously restricts the types of junctions allowable.

#### **Example of Parasitic Mode Excitation**



A slot-line mode is excited at a CPW junction

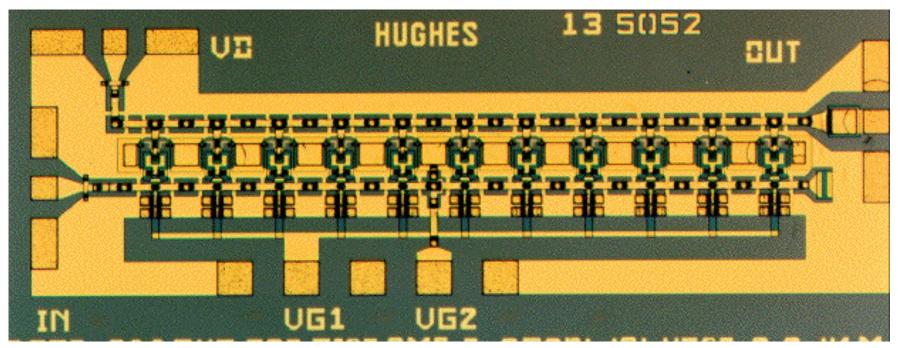


The fix...

this is one of many possible examples...

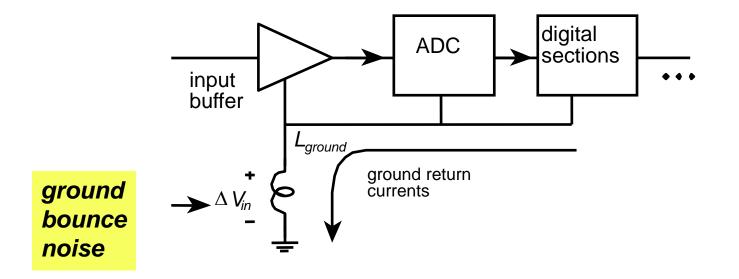
#### **Example of IC using CPW wiring**

1-180 GHz HEMT amplifier (UCSB / HRL) Note the ground bridges



# package resonance and grounding

#### What is Ground Bounce ?

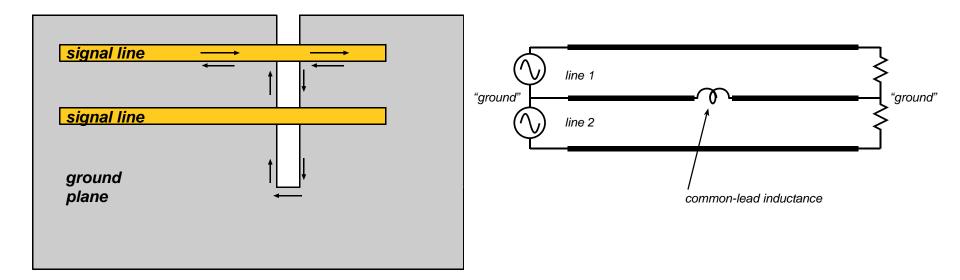


"Ground" simply means a reference potential shared between many circuit paths.

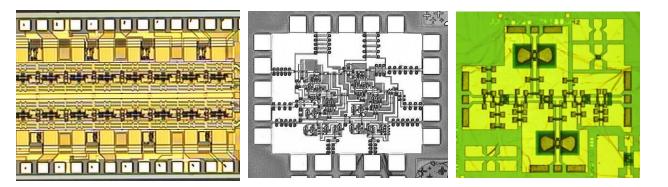
To the extent that it has nonzero impedance, circuits will couple in unexpected ways

RFI, resonance, oscillation, frequently result from poor ground systems

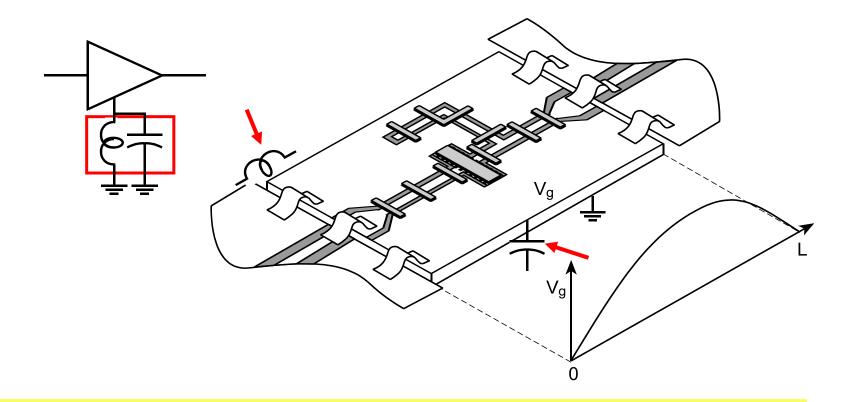
### Ground Bounce on an IC: break in a ground plane



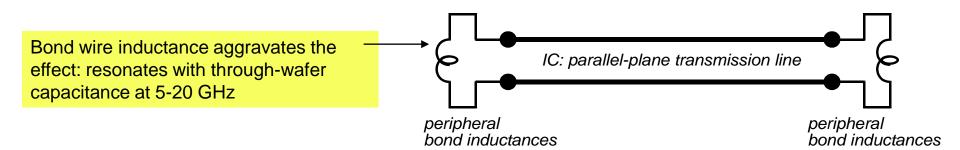
coupling / EMI due to poor ground system integrity is common in high-frequency systems whether on PC boards ...or on ICs.



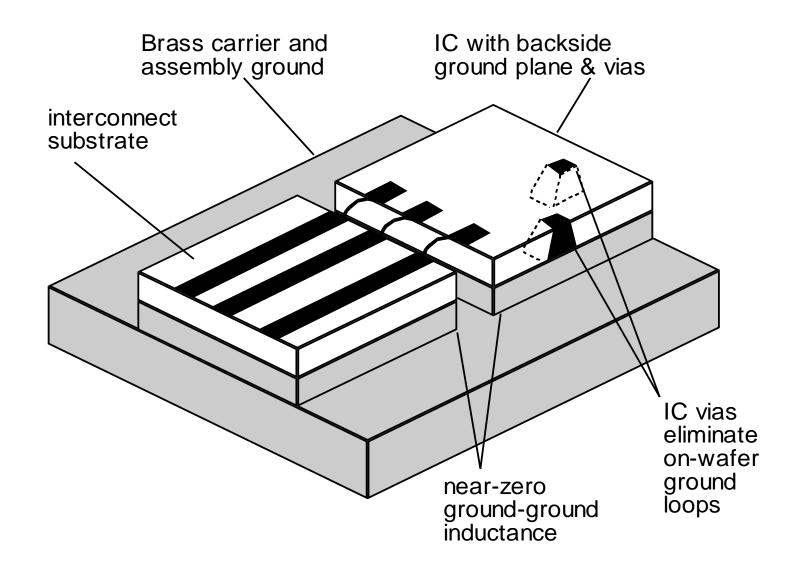
#### Ground Bounce: IC Packaging with Top-Surface-Only Ground



Peripheral grounding allows parallel plate mode resonance die dimensions must be <0.4mm at 100GHz



#### **Substrate Microstrip: Eliminates Ground Return Problems**

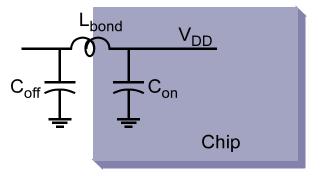


# power-supply resonance

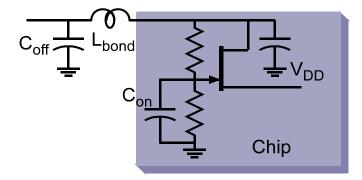
#### **Power Supply Resonance**

Resonates at  $f = 1/2\pi \sqrt{L_{bond}C_{on}}$ 

gain peak / suckout, oscillation, etc.



Active (AC) supply regulation

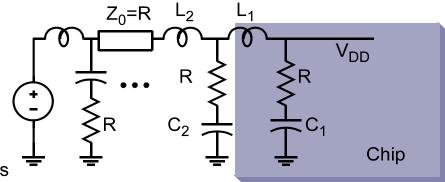


Passive filter synthesis

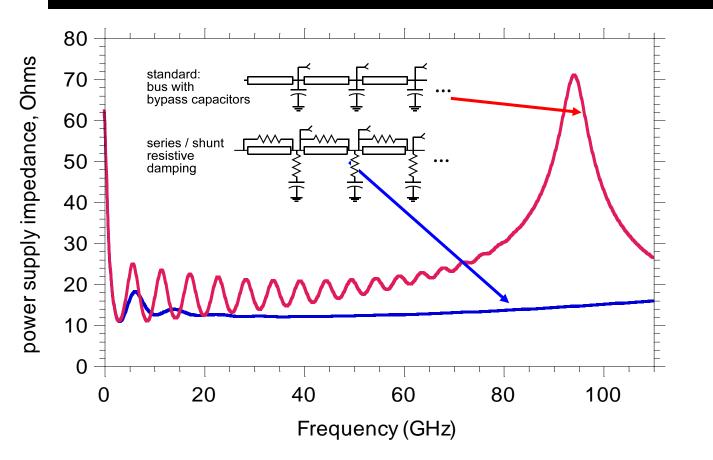
$$R = \sqrt{L_1/C_1}$$

$$\sqrt{L_1/C_1} = \sqrt{L_2/C_2} = \cdots$$

supply impedance is R at all frequencies



#### **Power Supply Resonances; Power Supply Damping**

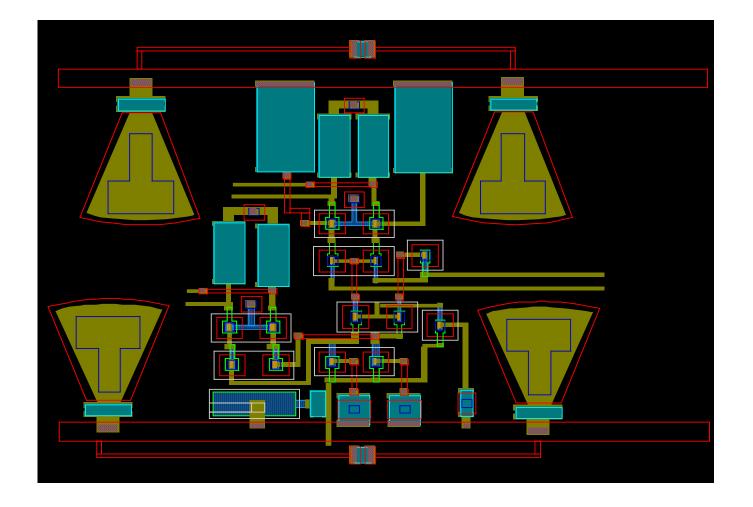


90 GHz--local resonance between power supply capacitance and supply lead inductance

~N\*5GHz resonances--global standing wave on power supply bus

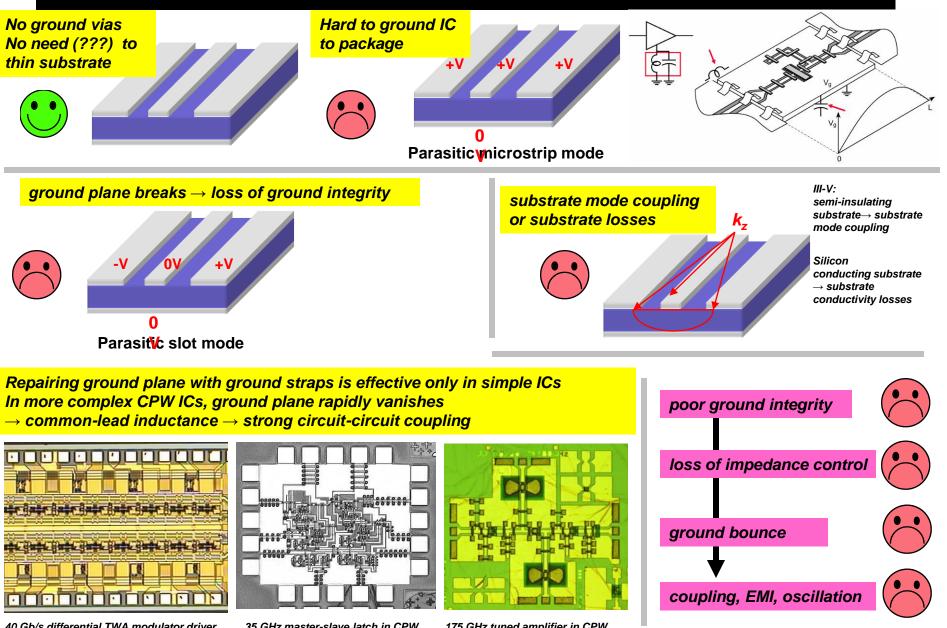
Power supply is certain to resonate: we must model, simulate, and add damping during design.

#### **Standard cell showing power busses**



# Interconnects: Summary, Design Strategy

# **Coplanar Waveguide: Summary**

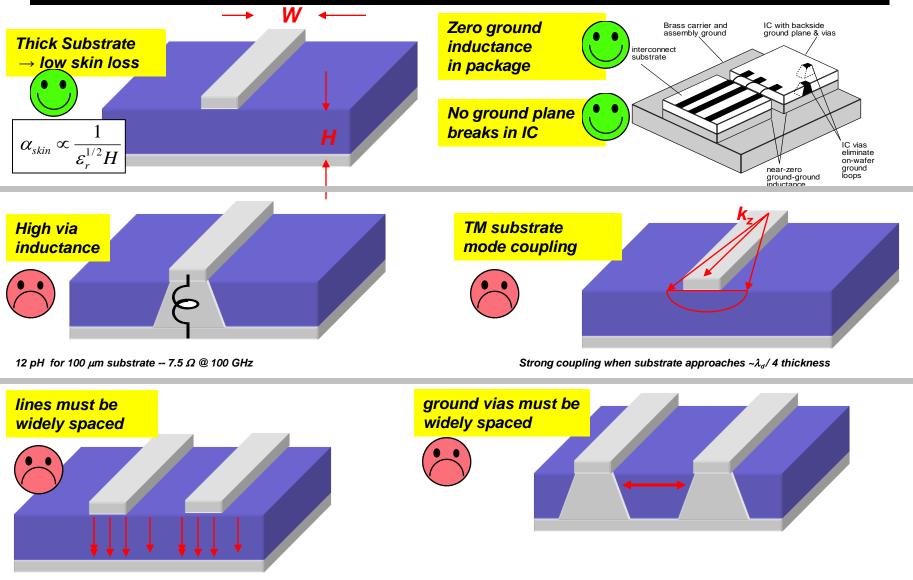


40 Gb/s differential TWA modulator driver note CPW lines, fragmented ground plane

35 GHz master-slave latch in CPW note fragmented ground plane

175 GHz tuned amplifier in CPW note fragmented ground plane

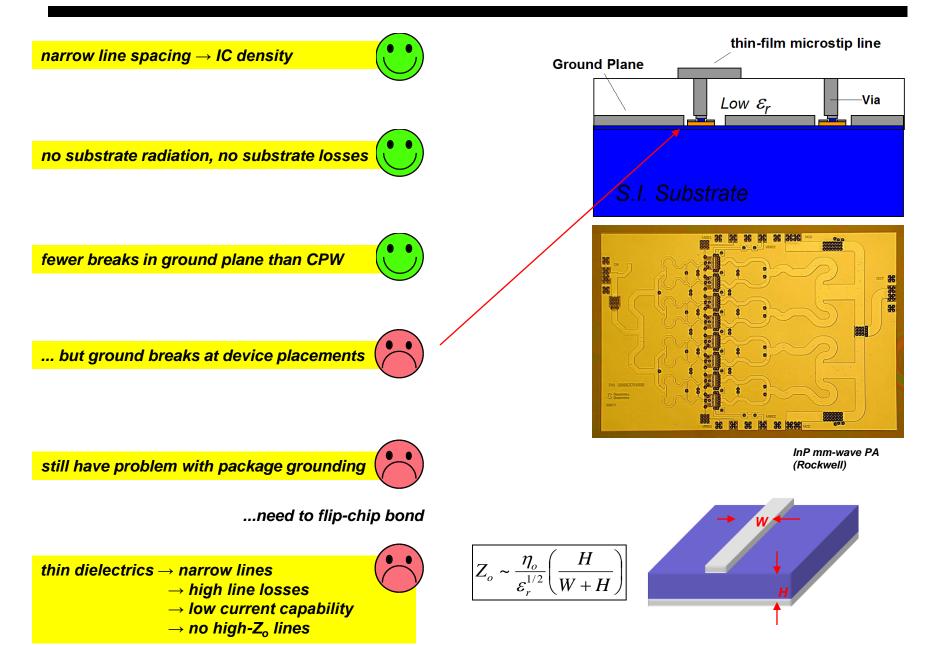
### **Classic Substrate Microstrip: Summary**



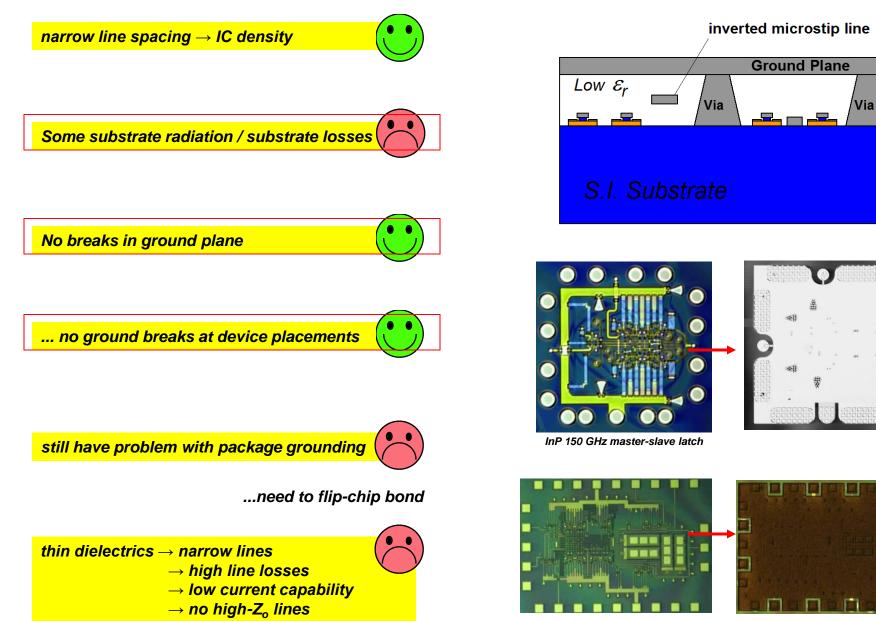
Line spacings must be ~3\*(substrate thickness)

all factors require very thin substrates for >100 GHz ICs  $\rightarrow$  lapping to ~50  $\mu$ m substrate thickness typical for 100+ GHz

## III-V MIMIC Interconnects -- Thin-Film Microstrip

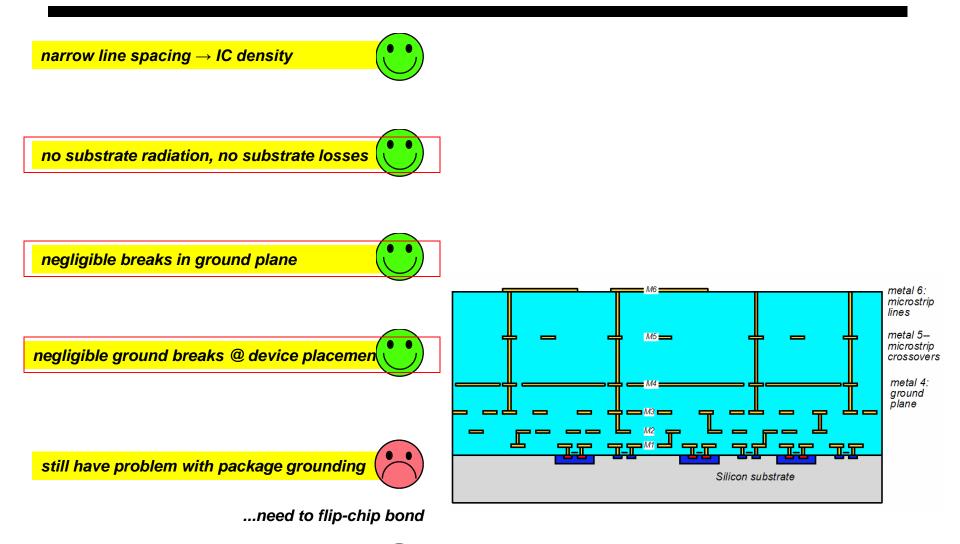


### III-V MIMIC Interconnects -- Inverted Thin-Film Microstrip



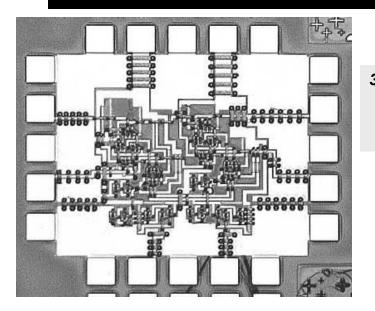
InP 8 GHz clock rate delta-sigma ADC

## VLSI Interconnects with Ground Integrity & Controlled Z<sub>o</sub>

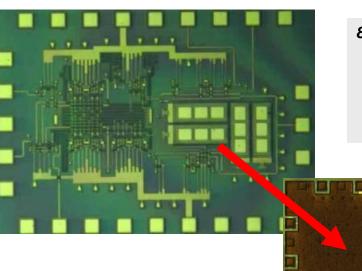


 $\begin{array}{l} \textit{thin dielectrics} \rightarrow \textit{narrow lines} \\ \rightarrow \textit{high line losses} \\ \rightarrow \textit{low current capability} \\ \rightarrow \textit{no high-Z}_{o} \textit{lines} \end{array}$ 

#### No clean ground return ? $\rightarrow$ interconnects can't be modeled !



35 GHz static divider interconnects have no clear local ground return interconnect inductance is non-local interconnect inductance has no compact model



8 GHz clock-rate delta-sigma ADC thin-film microstrip wiring every interconnect can be modeled as microstrip some interconnects are terminated in their Zo some interconnects are not terminated ...but ALL are precisely modeled

InP 8 GHz clock rate delta-sigma ADC

# End

# Appendix (optional)

## Skin effect losses l

Given a plane wave perpendicularly incident in direction z onto a sheet of metal :

$$\frac{\partial E}{\partial z} = -j\omega\mu H$$
 and  $\frac{\partial H}{\partial z} = -(j\omega\varepsilon + \sigma)E$ 

Hence  $E(z) = E_o e^{-\gamma z}$  where  $\gamma = \sqrt{j\omega\mu(j\omega\varepsilon + \sigma)}$ 

If 
$$\omega \varepsilon \ll \sigma$$
, then  $\gamma \cong \sqrt{j\omega\mu\sigma} = \sqrt{\omega\mu\sigma/2} + j\sqrt{\omega\mu\sigma/2}$ 

Defining the skin depth as  $\delta = \sqrt{2/\omega\mu\sigma}$  we find that  $E(z) = E_o e^{-z/\delta} e^{-jz/\delta}$ 

... the field dies down exponentially with distance into the metal.

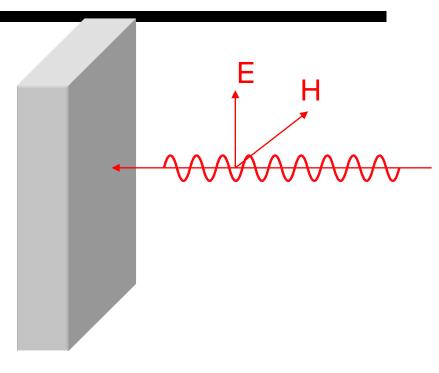
Wave impedance in the metal is :

$$\eta_{metal} = \frac{E}{H} = \sqrt{\frac{j\omega\mu}{j\omega\varepsilon + \sigma}} \cong \sqrt{\frac{j\omega\mu}{\sigma}} = \sqrt{\frac{\omega\mu}{2\sigma}} + j\sqrt{\frac{\omega\mu}{2\sigma}}$$

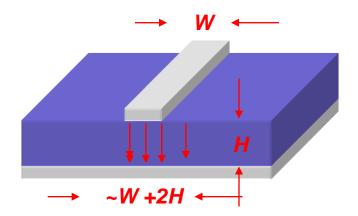
hence,

 $\eta_{metal} = \frac{1}{\sigma\delta} + j\frac{1}{\sigma\delta} \leftarrow \text{note the resistive and inductive terms}$ 

This is the SURFACE IMPEDANCE.



## Skin effect losses II



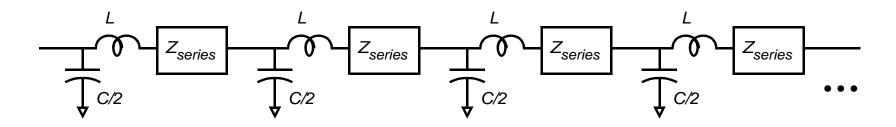
In a transmission line, the wave travels parallel, not perpendicular to the metal surface, but the same surface impedance is seen, provided that the transmission - line wavelength is much larger than the skin depth

The transmission - line then has an added series impedance per unit distance of  $Z_{series} = \frac{1}{P} \frac{(1+j)}{\delta \sigma}$ , where P is the effective current - carrying periphery.

For this microstrip line, there is surface impedance both in the signal and ground lines

$$Z_{series} \approx \left(\frac{1}{W} + \frac{1}{W + 2H}\right) \frac{(1+j)}{\delta\sigma}$$

## **Skin effect losses III**



This then introduces both loss and dispersion

$$Z_{series} = \frac{1}{P} \frac{(1+j)}{\delta \sigma} \quad \rightarrow \frac{\partial V}{\partial z} = -(j\omega L + Z_{series})I \text{ and } \frac{\partial I}{\partial z} = -j\omega CV$$

$$\rightarrow Z_o = \frac{V^+(z)}{I^+(z)} = \sqrt{\frac{j\omega L + Z_{series}}{j\omega C}} = \sqrt{\frac{j\omega L + 1/P\delta\sigma + j/P\delta\sigma}{j\omega C}}$$

....some secondary change in characteri stic impedance

$$\rightarrow V(z) = V_o e^{-\gamma_{line} z}, \text{ where}$$

$$\gamma_{line} = \sqrt{(j\omega L + Z_{series}) j\omega C} = j\omega \sqrt{LC} \sqrt{1 + Z_{series} / j\omega L}$$

$$\cong j\omega \sqrt{LC} (1 + Z_{series} / j2\omega L) = j\omega \sqrt{LC} + Z_{series} \left(\sqrt{C/L}\right) / 2$$

$$\gamma_{line} = j\omega \sqrt{LC} + Z_{series} / 2Z_o = j\omega \sqrt{LC} + \frac{1}{2Z_0 P \delta \sigma} + \frac{j}{2Z_0 P \delta \sigma}$$
Skin Loss dispersion

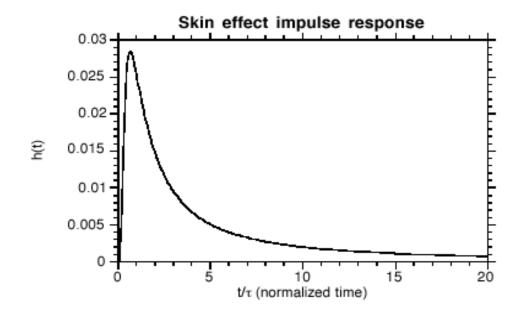
## Skin Effect losses, IV

The impulse response of the transmission line can then be found. (Wiginton and Nahman, Proc. IRE, February 1957)

$$h(t) \approx C * U(t/\tau) \frac{(t/\tau)^{-3/2} \exp(-\tau/t)}{\tau \sqrt{\pi}}$$

where 
$$\tau = \left[ l_{\sqrt{\frac{\mu}{\sigma}}} / 4Z_0 P \right]^2$$

Skin effect causes pulse broadening proportional to distance<sup>2</sup>



## **Skin effect losses V**

The step response is the integral of the impulse response. Note the initial fast rise and the subsequent "dribble-up" characteristic of skin effect losses.

