

# • Notes Set 23: Fiber optic communications systems

- Receiver design in colored noise. Whitening filter. Intractability of problem
- sketch of decision-feedback strategy.
- Simplified practical (nonoptimum) receivers. Eye diagrams. Personik noise integrals. Optical receiver sensitivity
- design of low-noise optical receiver front ends

①

FCE — , Notes set 23

Fiber optic digital Communications:

- How it differs from prior materials
- Smith & Persenik's Method
- Performance Estimates.

We need to return to our prior material:

optimum receiver strategy:

1) Message sent by message vector / sequence  
- represented by basis functions  $\phi_i$

2) Message now usually detected against  
white Gaussian noise.

3) Project signal against basis vectors.

then project signal vector against set  
of possible signals, to see which  
is closest.

what happens if noise is not white?

ex 11  $\star$  narrowband digital microwave receiver

$$F = \text{noise figure} = F_0 = \text{constant across band.}$$

$\Rightarrow$  white system noise

$\star$  Very broadband digital microwave receiver:

Fano's limit: cut math for noise across whole bandwidth

$$F = F_0 + K(f - f_0)^2$$

↑  
channel center frequency.

- noise no longer white.

$\star$  Broadband PCM fiber optics

$$\left. \begin{array}{l} \text{Input noise current} \\ \text{spectral density} \end{array} \right\} = K_1 + K_2 f^2$$

$\Rightarrow$  Not white

What does the theory tell us to do

\* Pass the signal + noise through filter

such that output noise spectrum is white

"whitening filter"

\* Theory of previous lecture

- in theory -

... gives us the correct receiver architecture.

Problem:

\* After whitening Filter,

the message vectors are, bit-by-bit  
no longer orthogonal

\* Theory acknowledges no problem:

architecture given previously is "optimum"

\* The problem: the "optimum" receiver

structure is indeed optimum in terms of  
receiver sensitivity, but is highly  
non-optimum in terms of cost & difficulty  
of implementation.

Example to clarify this:

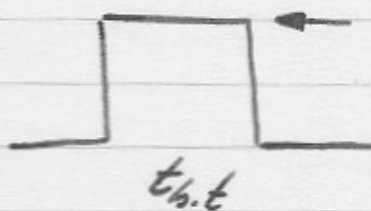
Fiber-optic receiver:

\* input referred noise power spectral density

$$N_0 \approx a + bf^2$$

\* Signalling Format: on-off Keying.

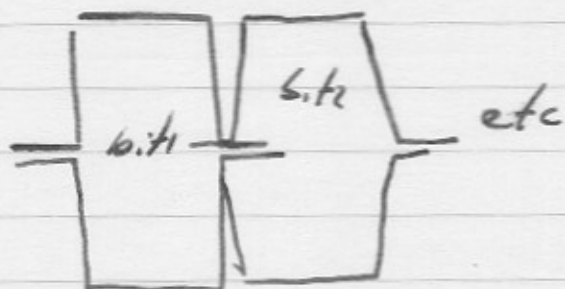
\* Basis Function



The Fiber people call this Non-Return-to-zero (NRZ)

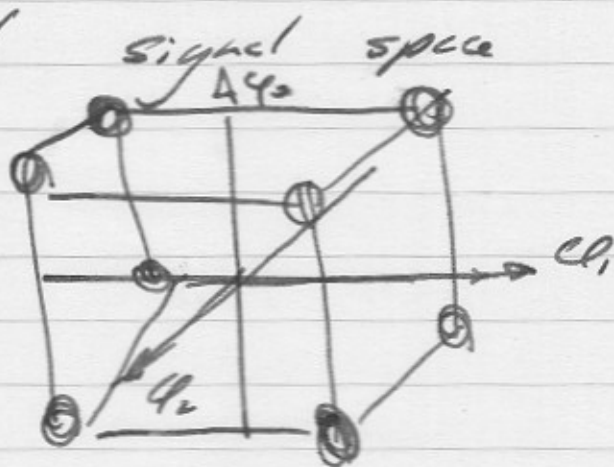
Better Name: OOK with square symbol.

Note that the bits are orthogonal



... so that, although we could in principle look at  $N$  bits as  $N$  boundaries in an

$N$ -dimensional

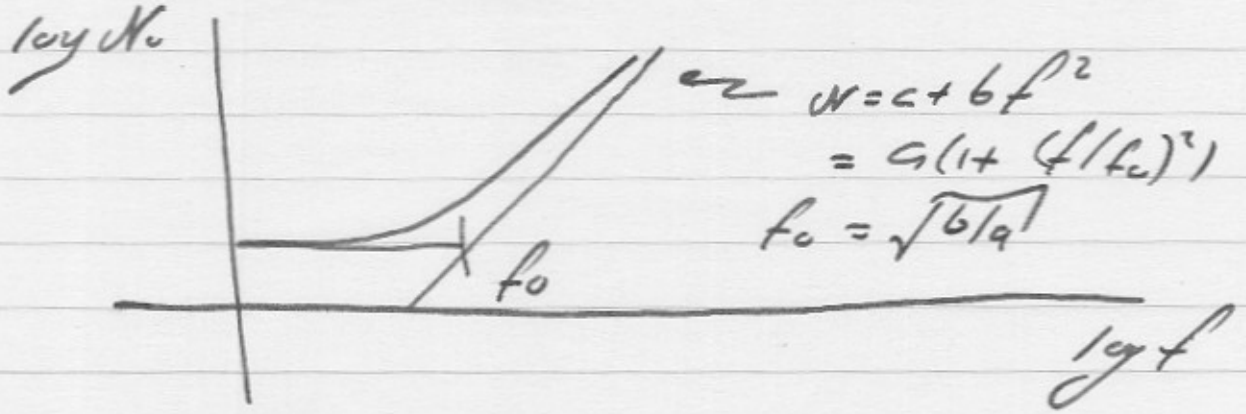


we can in fact look at one bit at a time

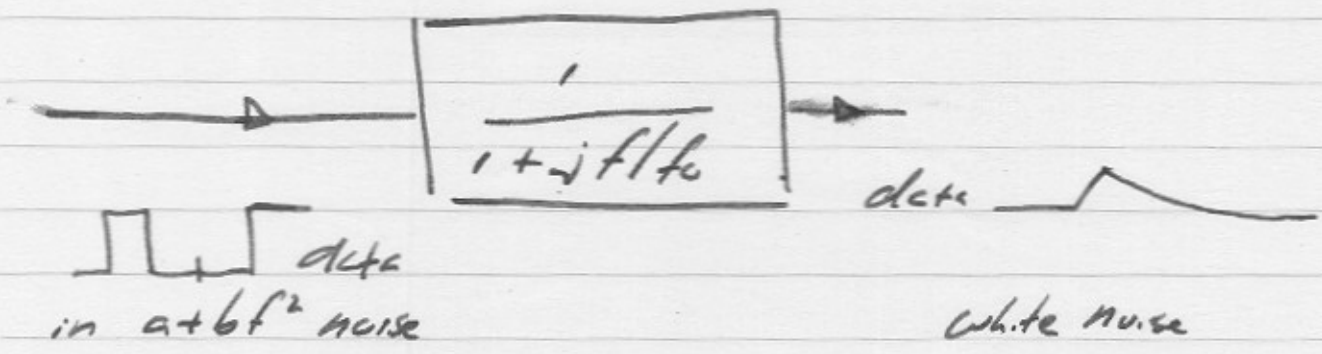
↳ nice & simple receiver



\* but the noise has p.s.d. like so:



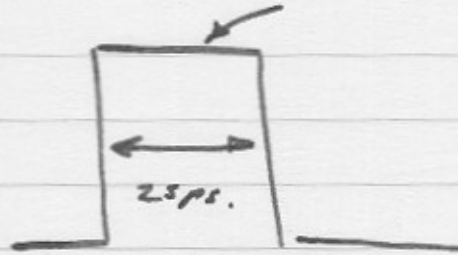
So an optimal receiver must pass the signal + noise first through a whitening filter:



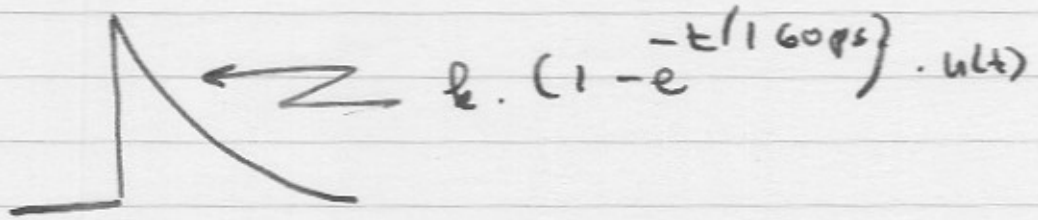
TBL 40 Gbps system:  $T_b = 25 \text{ ps}$

Bipolar receiver:  $f_0 \sim 1 \text{ GHz}$  ( $\tau_0 = 160 \text{ ps}$ )

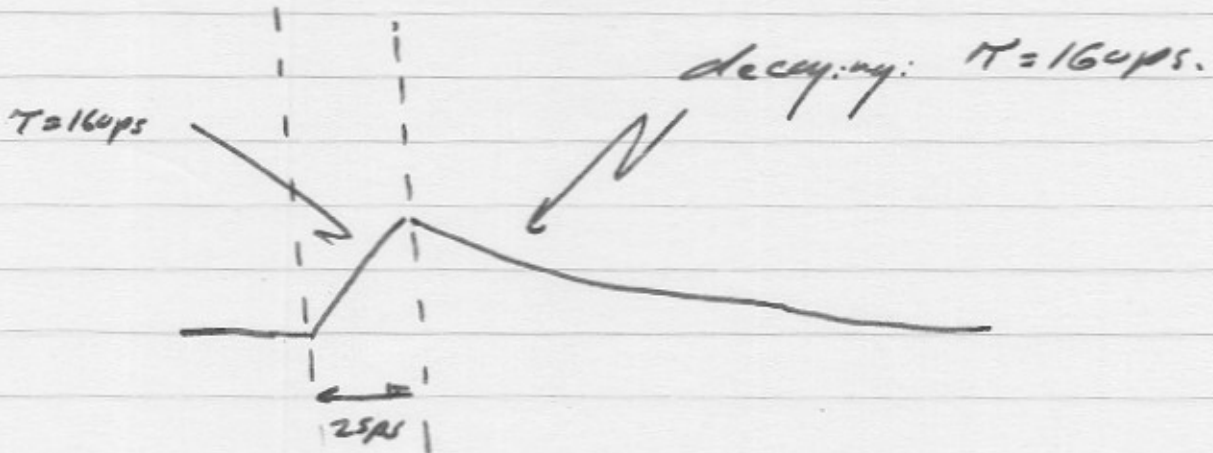
Input Pulse to filter:



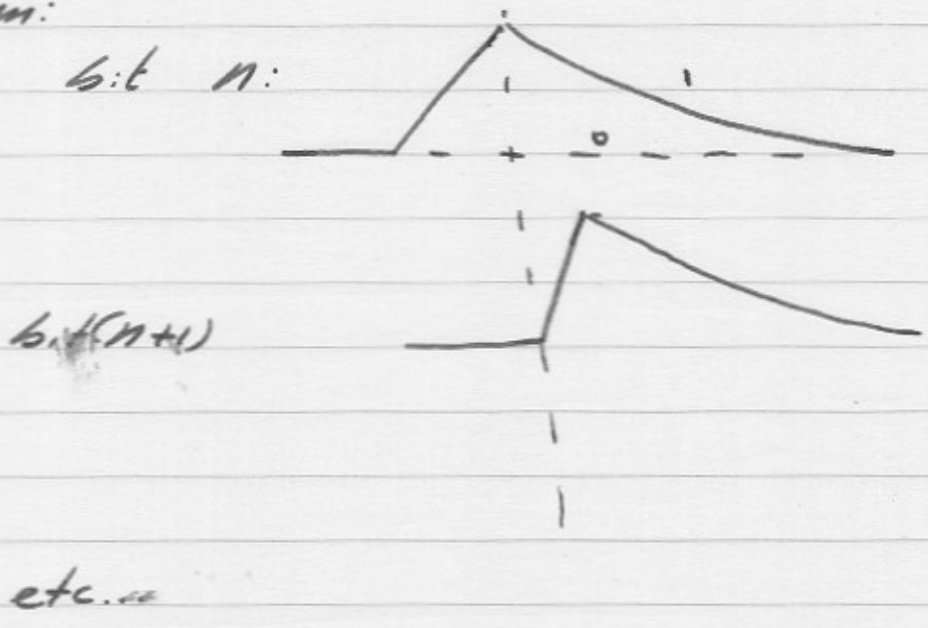
output  
Impulse response of filter:



output pulse from filter



problem:

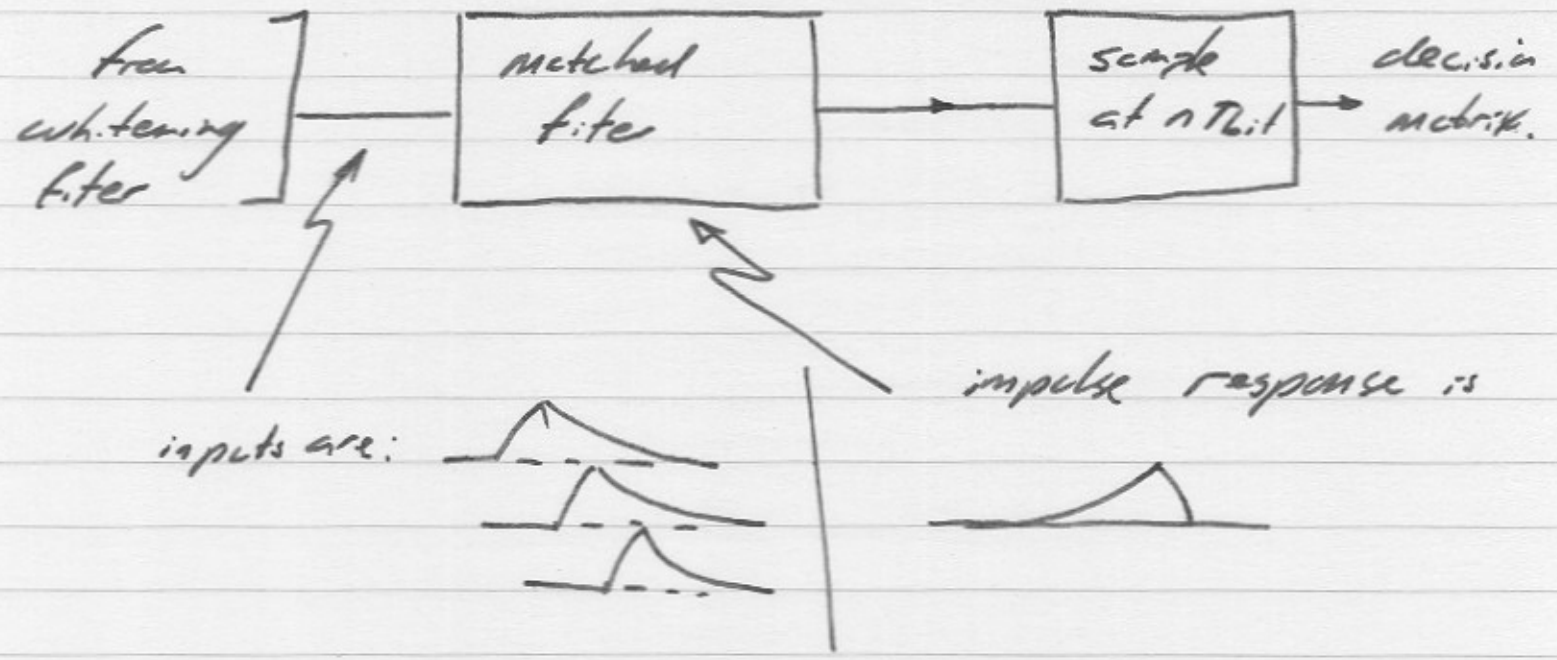


etc...

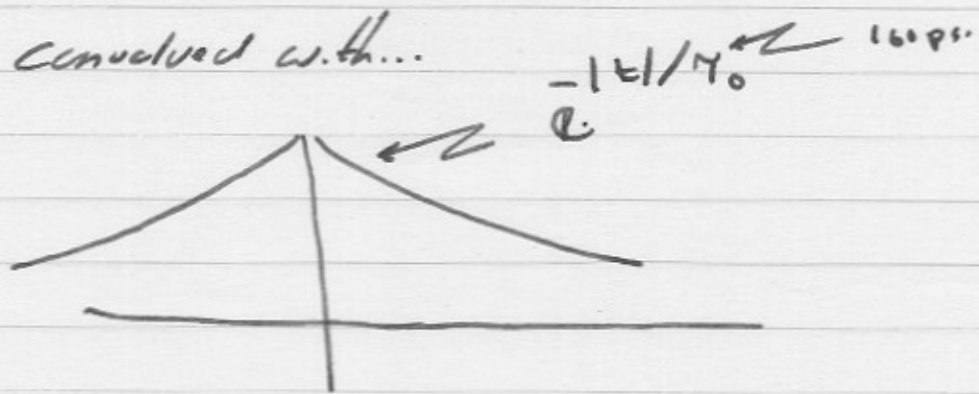
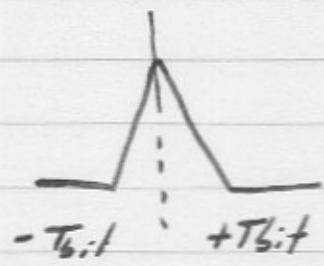
→ data bits no longer orthogonal.

→ our theory did not demand orthogonal signals to give an optimum receiver, but it did demand orthogonal bits if that optimum receiver were to be simple

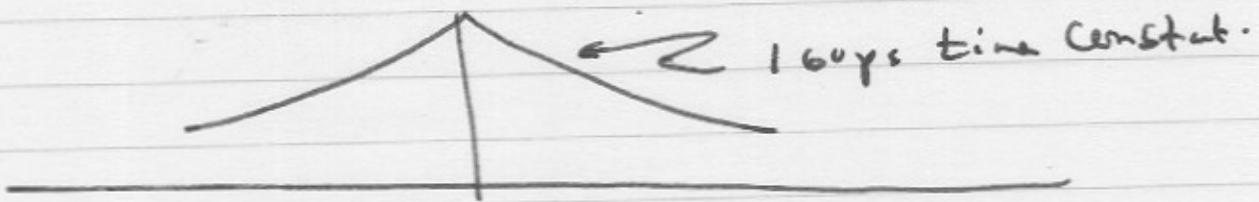
# Receiver structure "optimum"



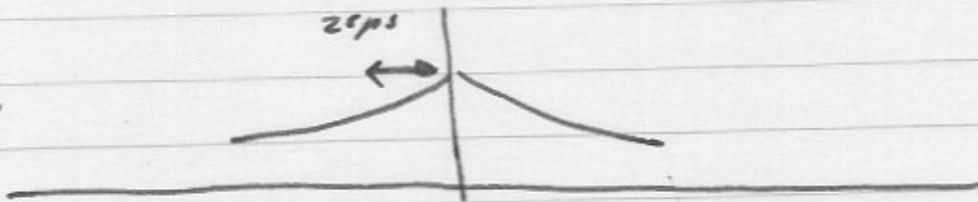
→ outputs of matched filter is therefore:



Response to single bit is  $\approx$  approximately:



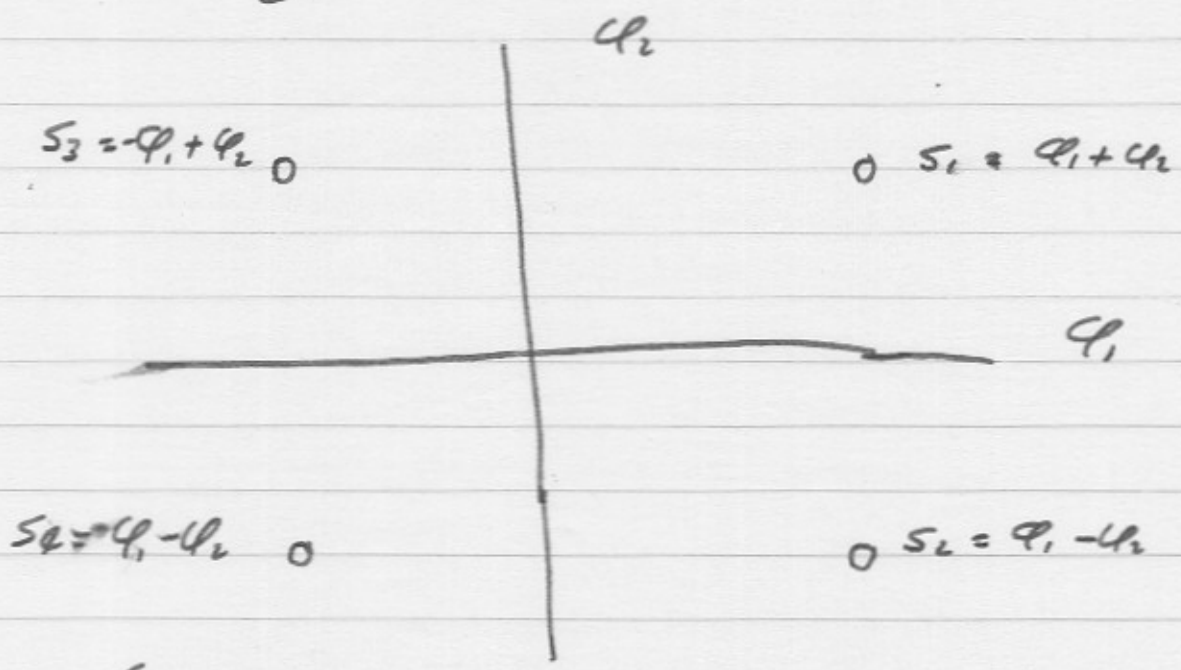
Next bit



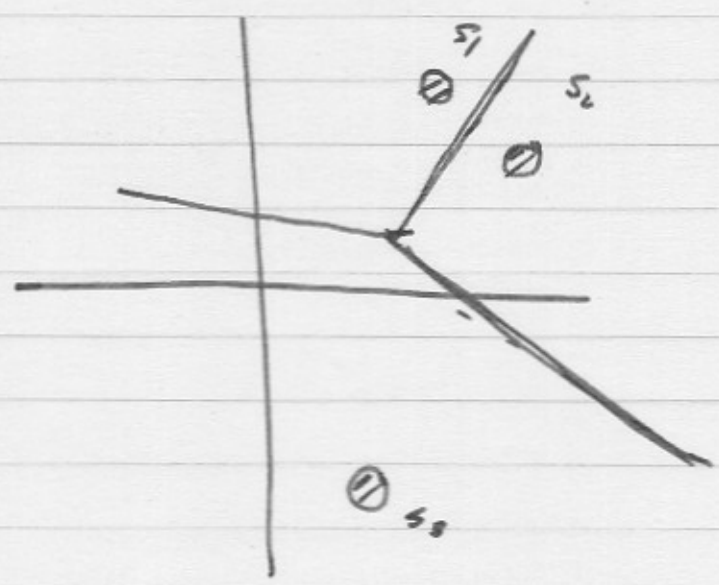
etc:

Very strong intersymbol interference.

again, Intersymbol Interference is ok:



no intersymbol interference allows us to decide on bit-by-bit basis ( $q_1, q_2, q_3, \dots$ )



these pictures pertain to the non-orthogonal signal case

... but with 25ps-duration bits spread  
 in time to  $\pm 160\text{ps}$ , the decision algorithm  
 must consider  $\sim 2 \cdot (160/25) = 13$  successive  
 outputs of the correlator / matched filter in  
 order to determine a given bit's value.

↓  
 Ideal hardware becomes strictly  
 unrealizable

⇓  
 Sacrifice some performance for some reduction  
 in complexity: decision feedback method.

- or -

sacrifice substantial performance for  
 extreme hardware simplification  $\rightarrow$  Nyquist Filter Method.

\* I will not discuss  
decision Feedback here.

\* Nyquist Method:

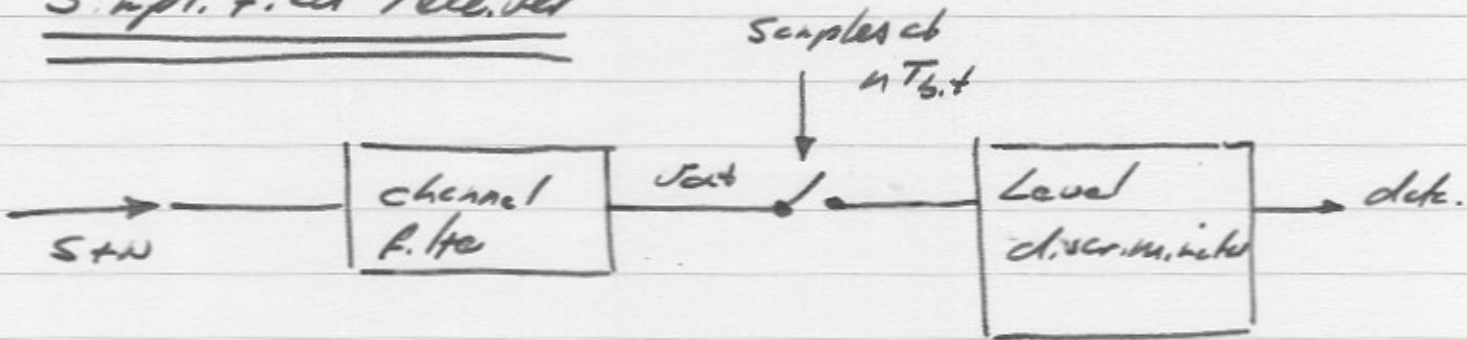
1. Extremely simple hardware.

2. Replace whitening & matched filter/correlator  
with simple noise-bandlimiting filter

\* subject to constraint, zero  
intersymbol interference: E.G.

Orthogonal bit slots



Simpl. Fied receiver

(1-0-0)

\* Channel Ltr constraint:

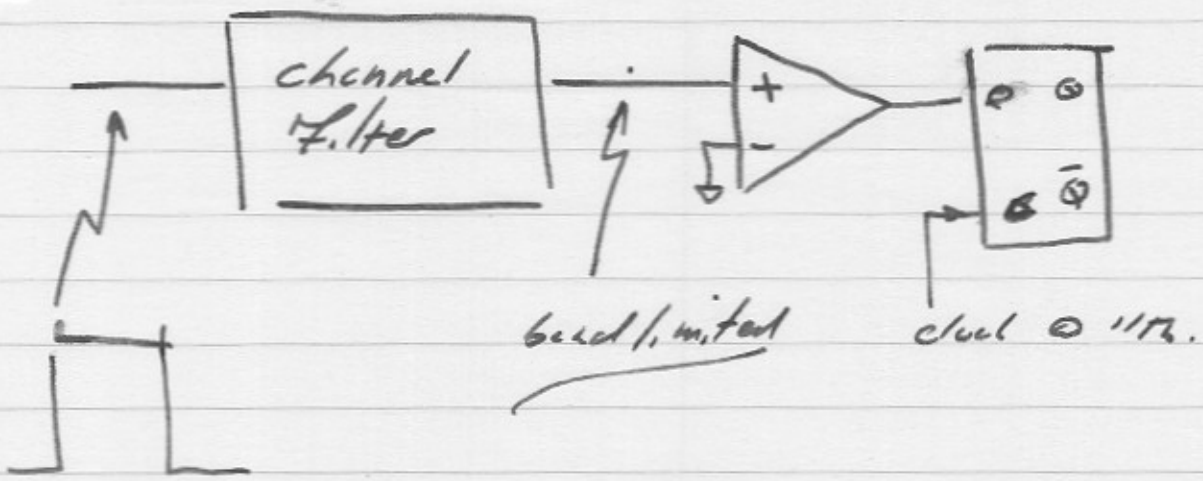
$\equiv V_{out}(nT_{b,t})$  depends only on value of  $n^{\text{th}}$  bit  
 - zero ISI - intersymbol interference.

$\equiv$  try to minimize noise...

\* order of { sampler } and d. scr. min. ltr

con & often is reversed  $\rightarrow$  becomes

comparator & clocked latch.

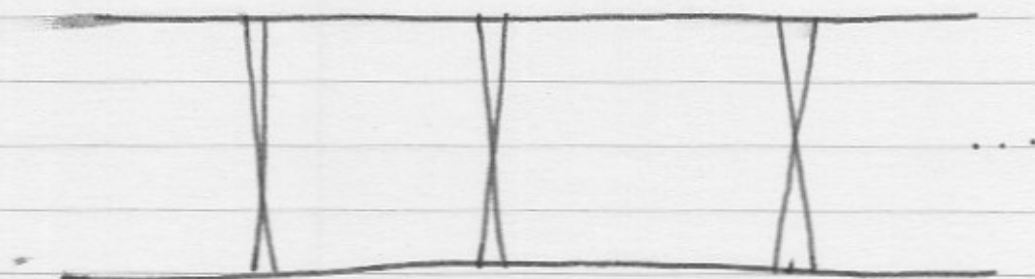


# Eye diagram

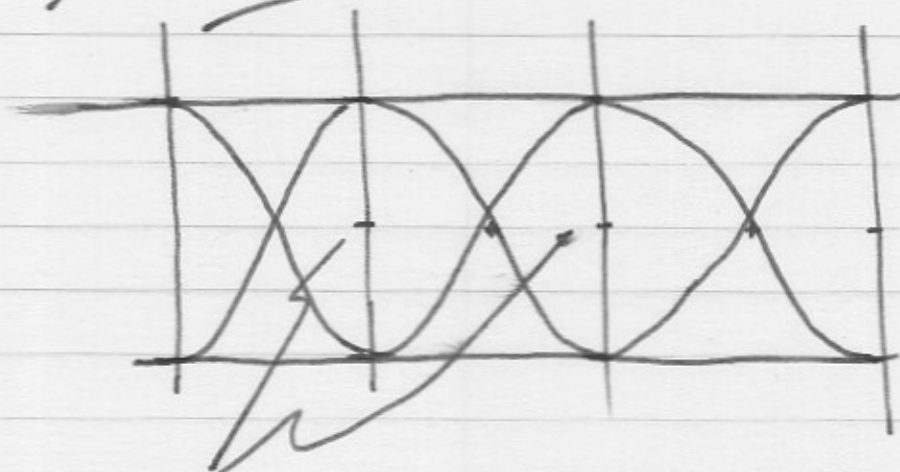
Born on an oscilloscope:

- scope triggered on clock
- scope displays data stream
- all possible trajectories superimposed.

## Input Eye

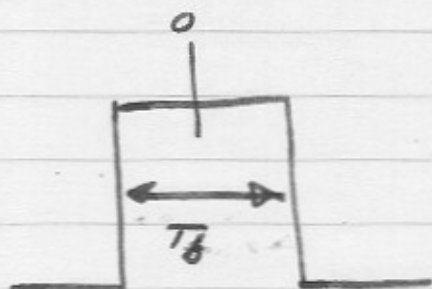


## Output Eye (a good one)

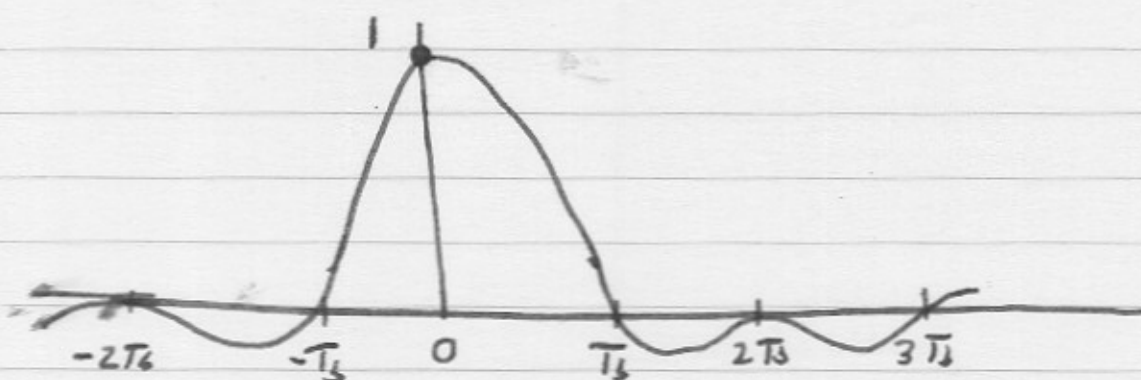


decision circuit looks at voltage at these points in time...

This demands that a bit input:



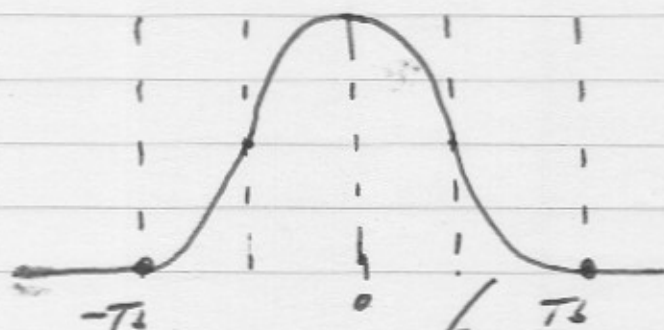
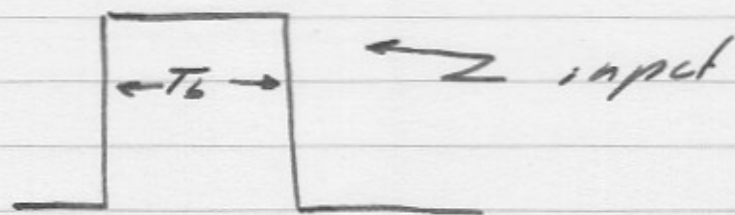
Produce (models a time delay) an output:



... which is zero at the center of all other bits.

Filtering filters that do this is a general problem solved by Nyquist they are called Nyquist filters, although not pertaining to Sampling

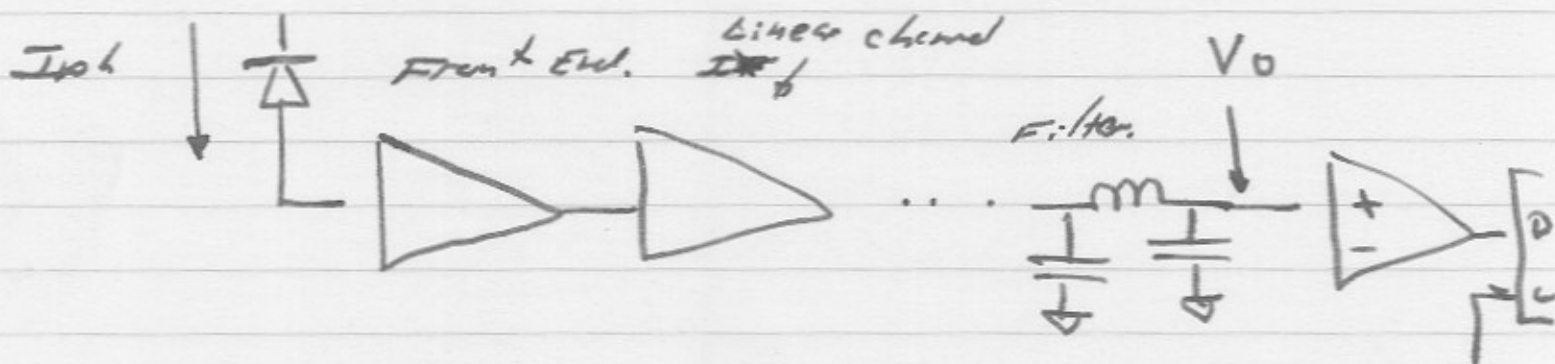
Smith & Peryonik considered  
a far more restrictive set of (idealized,  
not real) filters



raised cosine:

$$\frac{\sin(\pi t/T_b) \cos(\pi t/T_b)}{(\pi t/T_b) (1 - (2t/T_b)^2)}$$

How they do the analysis...



$$I_{ph} = P_{optical} \cdot \frac{q}{h\nu} \cdot \text{detector quantum efficiency.}$$

$\swarrow$   
 $\swarrow$   
 $\swarrow$  1 electron per photon

$$V_o = K Z_T(\omega) I_{ph}$$

$\swarrow$   
 pick K so that  $Z_T(\omega) \rightarrow 1$  at DC.

\* We analyze input network to find total input-referred noise current:

$$\frac{d \langle i_{in}^2 \rangle}{df} = a + bf^2$$



more on this in a minute...

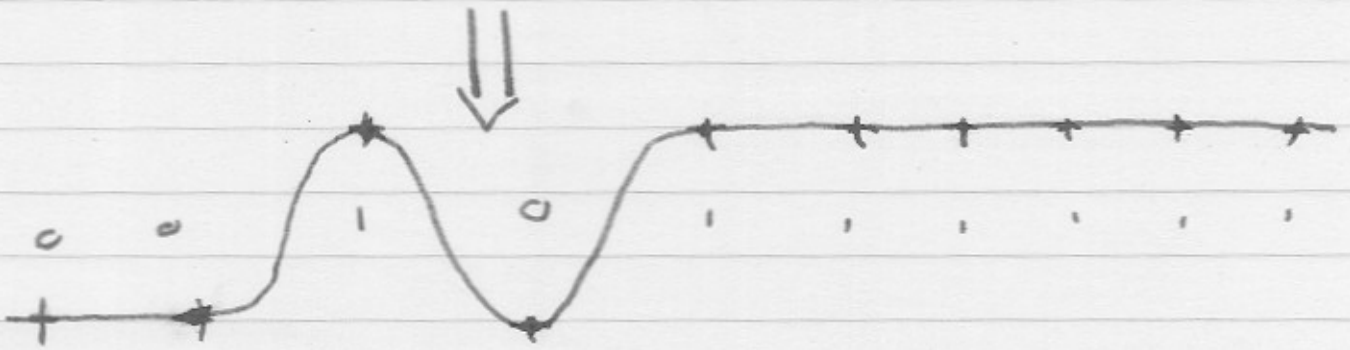
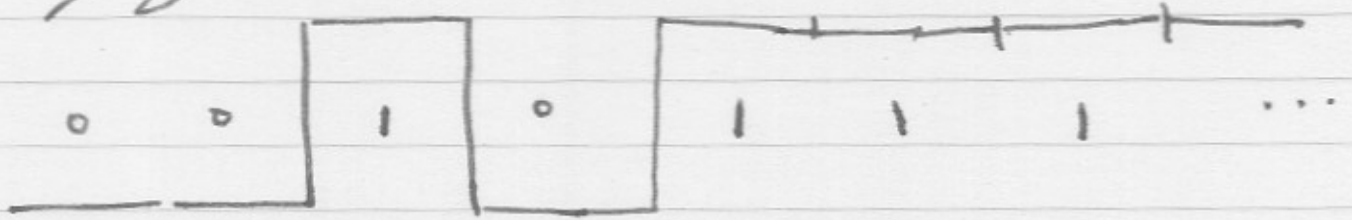
\* The output noise voltage has variance

$$\sigma_{V_{in}}^2 = k \int_0^{+\infty} \|Z_T(\omega)\|^2 \cdot (a + bf^2) \cdot df$$

\* The output signal has magnitude:

$$V_{out} = I_{in} \cdot K Z_T(\omega=0)$$

\* This has the key assumption of  
a Nyquist zero-ISZ channel filter:



... so that the amplitude of

an isolated bit is indeed  $I_m \cdot k_T(\omega_c)$

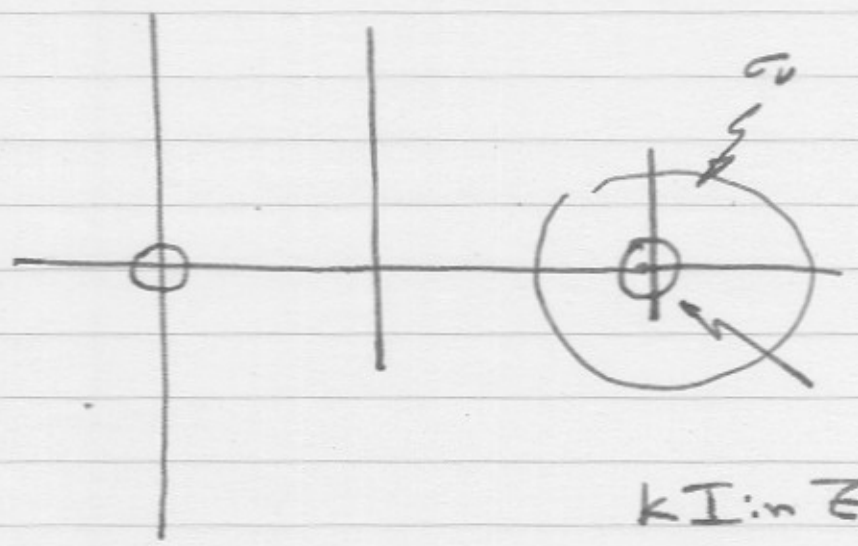


Then  $\sigma_{vn} = K \sqrt{\int_0^\infty \|Z_T(f)\|^2 (c + bf^2) df}$

and  $V_{cut} = K I_{in} Z(0)$

OOK

boundary  $K I_{in} Z(0)/2$



write  $\sigma_{vn} = K X$  for a moment.

$$P(\text{Error}) = Q \left[ \frac{(k I_{in} Z(0) / 2)}{kX} \right]$$

$$= Q \left[ \frac{I_{in} Z(0) / 2X}{1} \right]$$

as usual,  $P(\text{Error}) = \text{e.g.} = 10^{-9}$

$$\rightarrow \frac{I_{in} Z(0) / 2X}{1} = "Q" = 6$$

$$\frac{I_{in}}{Z(0)} = \frac{2X \cdot Q}{Z(0)} = 2Q \sqrt{\int_0^{\infty} \left\| \frac{Z_T(f)}{Z_T(0)} \right\|^2 (a + bf^2) df}$$

where we have assumed  $S_I(f) = a + bf^2$

so it is general:

$$\frac{I_{in}}{Z(0)} = 2Q \cdot \sqrt{\int_0^{\infty} \left\| \frac{Z_T(f)}{Z_T(0)} \right\|^2 S_I(f) df}$$

Receiver Sensitivity:

$$1 \text{ sent} \quad P_{\text{opt}} = P_1 \Rightarrow I_{i_2} = \frac{\eta q}{h\nu} P_1$$

$$0 \text{ sent} \quad P_{\text{opt}} = 0 \Rightarrow I_{i_2} = 0$$

$$\bar{P} = \text{average power} = P_1/2$$

$$= I_{i_2} \frac{h\nu}{\eta q} \cdot \frac{1}{2}$$

$$\bar{P}_{\text{mic}} = \frac{h\nu}{78} \cdot Q \int_0^{\infty} \left\| \frac{z_T(f)}{z_T(0)} \right\|^2 \cdot S_I(f) df$$

now usually:  $S_I(f) = a + bf^2$

$$\langle I^2 \rangle \triangleq \int_0^{\infty} \left\| \frac{z_T(f)}{z_T(0)} \right\|^2 S_I(f) df$$

$$= \int_0^{\infty} \left\| \frac{z_T(f)}{z_T(0)} \right\|^2 \cdot a \cdot df$$

$$+ \int_0^{\infty} \left\| \frac{z_T(f)}{z_T(0)} \right\|^2 \cdot bf^2 df$$

Let's normalize with respect to  
 Bit rate:  $B$  bits/second.

$B$  bits/second  $\rightarrow$  filter  $Z_T(f)$

$2B$  bits/second  $\rightarrow$  filter  $Z_T(f/2)$

\* we want to work with normalized  
 filter frequency responses; e.g. normalized  
 to the bit rate...

$$\tilde{Z}_T(f/B) = Z_T(f)$$

$$\langle I_1^2 \rangle = \int_0^\infty \left\| \frac{\tilde{Z}_T(f/B)}{\tilde{Z}_T(0)} \right\|^2 \cdot a \cdot df$$

$$+ \int_0^\infty \left\| \frac{\tilde{Z}_T(f/B)}{\tilde{Z}_T(0)} \right\|^2 \cdot b f^2 \cdot df$$

write  $\tilde{f} = f/B$

$$= aB \int_0^\infty \left\| \frac{\tilde{Z}_T(\tilde{f})}{\tilde{Z}_T(0)} \right\|^2 d\tilde{f}$$

$$+ bB^3 \int_0^\infty \left\| \frac{\tilde{Z}_T(\tilde{f})}{\tilde{Z}_T(0)} \right\|^2 \tilde{f}^2 d\tilde{f}$$

$$I_2 = \int_0^{\infty} \left\| \frac{\tilde{Z}_T(\tilde{f})}{\tilde{Z}_T(0)} \right\|^2 d\tilde{f}$$

$$I_3 = \int_0^{\infty} \left\| \frac{\tilde{Z}_T(\tilde{f})}{\tilde{Z}_T(0)} \right\|^2 \tilde{f}^2 d\tilde{f}$$

... are called Parzenik Coefficients.

They are the integrals of the filter's transfer functions in normalized frequency (normalized to the bit rate) and therefore depend only on the shape of the filtered pulse.

$$P_{min} = \frac{h\nu}{8} \cdot Q \cdot \frac{1}{\eta} \cdot \sqrt{\langle i_n^2 \rangle}$$

where

$$\langle i_n^2 \rangle = B \cdot I_2 \cdot a + B^3 I_3 \cdot b$$

$$\text{if } S_i(f) = a + bf^2$$

$$I_2 \approx 0.68$$

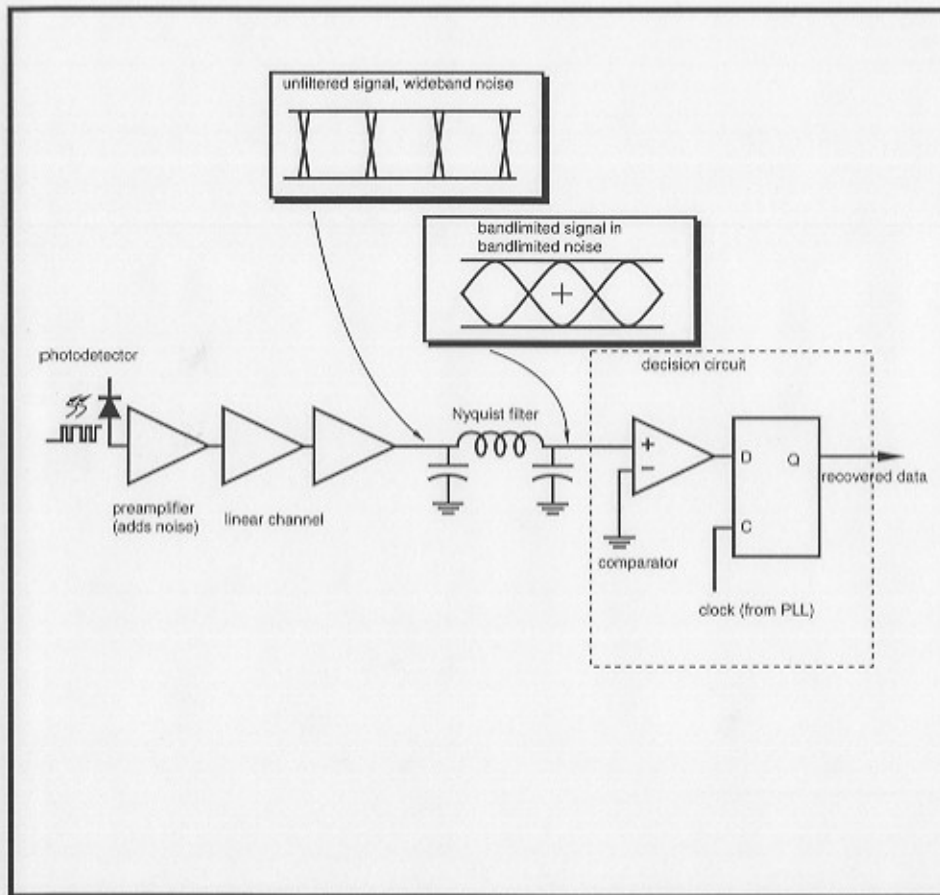
$$I_3 \approx 0.12$$

} parameters for a  
"real world" 5-pole L-C  
filter which approximates a raised  
cosine.

Raised cosine:  $I_2$  about the same  
 $I_3 \approx 0.06$



# Optical Receiver Basics



Photodetector converts light into current, drives preamplifier

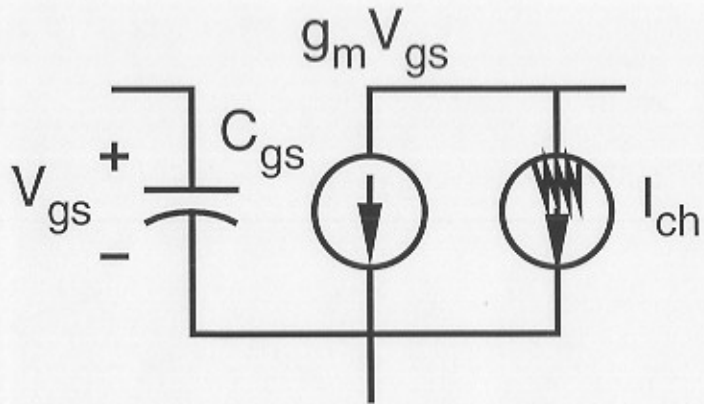
Preamplifier & photodetector set noise performance, hence attainable sensitivity

Nyquist filter bandlimits noise to  $\approx B/2$ .

Decision circuit recovers binary data

## Key Component Models

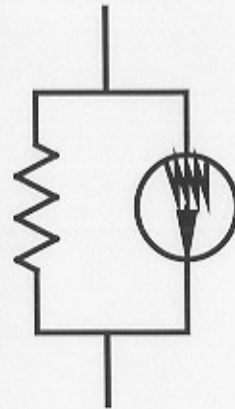
### Field-Effect Transistors



$$\frac{d}{df} \langle I_n I_n^* \rangle = 4kT \Gamma g_m$$

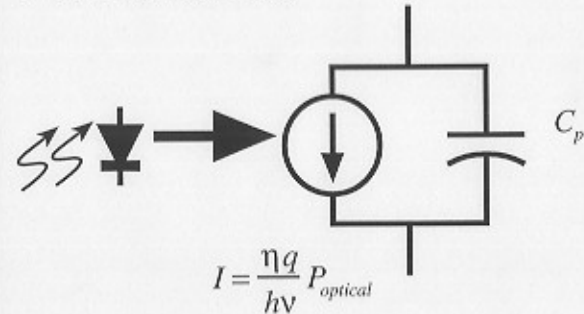
Other FET parasitics at  $kT$   
 1/f noise not relevant at rates above  
 10 GB/s

### Resistors



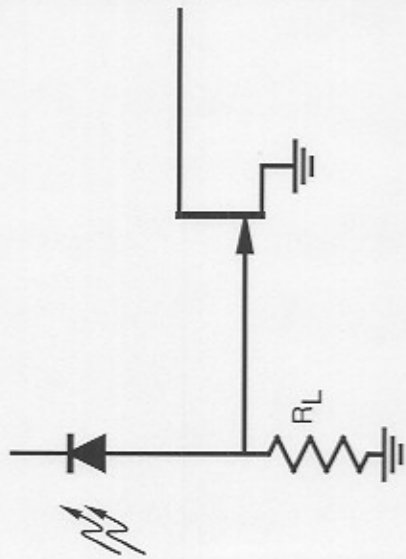
$$\frac{d}{df} \langle I_n I_n^* \rangle = 4kT / R$$

### Photodiode



$$I = \frac{\eta q}{h\nu} P_{\text{optical}}$$

## receiver noise analysis



Approximate input referred noise

$$\frac{d\langle I_n I_n^* \rangle}{df} = \frac{4kT}{R_L} + \frac{4kT\Gamma}{g_m} (2\pi f)^2 (C_{gs} + C_p)^2$$

but  $C_{gs} = g_m / 2\pi f_\tau$ , so choose  $C_{gs} = C_p$

hence:

$$\frac{d\langle I_n I_n^* \rangle}{df} = 32\pi kT\Gamma f^2 C_p / f_\tau + 4kT / R_f$$

## Front-End Noise and Bandwidth

The 'Simple' diagram shows a diode connected to a resistor  $R_L$ , which is then connected to the gate of a FET. The 'Transimpedance' diagram shows a similar setup but with an op-amp buffer connected to the FET gate, labeled 'gain=A'.

- Broadband circuit: cannot noise-match, noise does not approach  $f_{min}$ .
- Fundamental noise source is FET channel noise. Resistive loading added for bandwidth, degrades noise. Note: post-equalization often used.

$$\frac{d\langle I_n I_n^* \rangle}{df} \cong 32\pi kT \Gamma f^2 C_p / f_\tau + 4kT / R_f$$

This assumes intelligent choice of FET size

- Bandwidth is

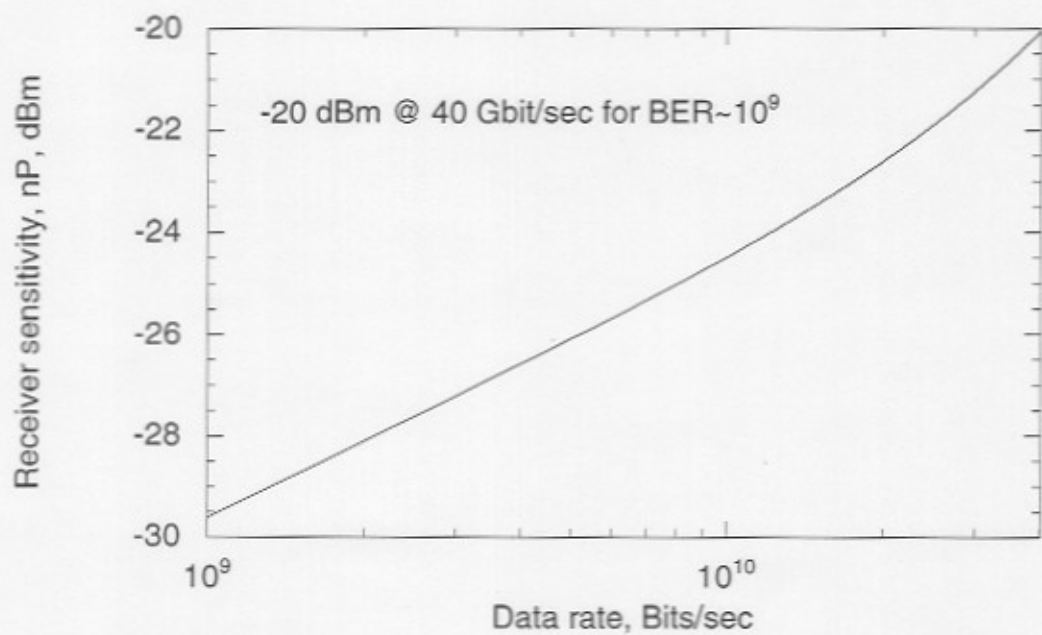
$$f_{3dB} \approx \frac{1 + A}{R_f (C_{diode} + C_{gs})}$$

(set  $A=1$  for simple stage)

## Common-Gate / Transimpedance Front-Ends

<p>same noise voltage and current in CS and CG modes.</p> <p>CG stage gives input pole at <math>f_\tau / 2</math>.</p> <p>need high <math>f_\tau</math>, small <math>C_{photodiode}</math>.</p>		
<p>input noise, given <math>C_{gs} = C_{photodiode}</math>:</p>	$\frac{d\langle I_n I_n^* \rangle}{df} \cong 32\pi kT \Gamma f^2 C_p / f_\tau + 4kT / R_f$	
<p>bandwidth:</p>	$f_{3dB} \approx \frac{1 + A}{R_f (C_{diode} + C_{gs})}$	$f_{3dB} \approx f_\tau / 2$

# Optical Receiver Sensitivity: Design Target



## Key observations regarding noise:

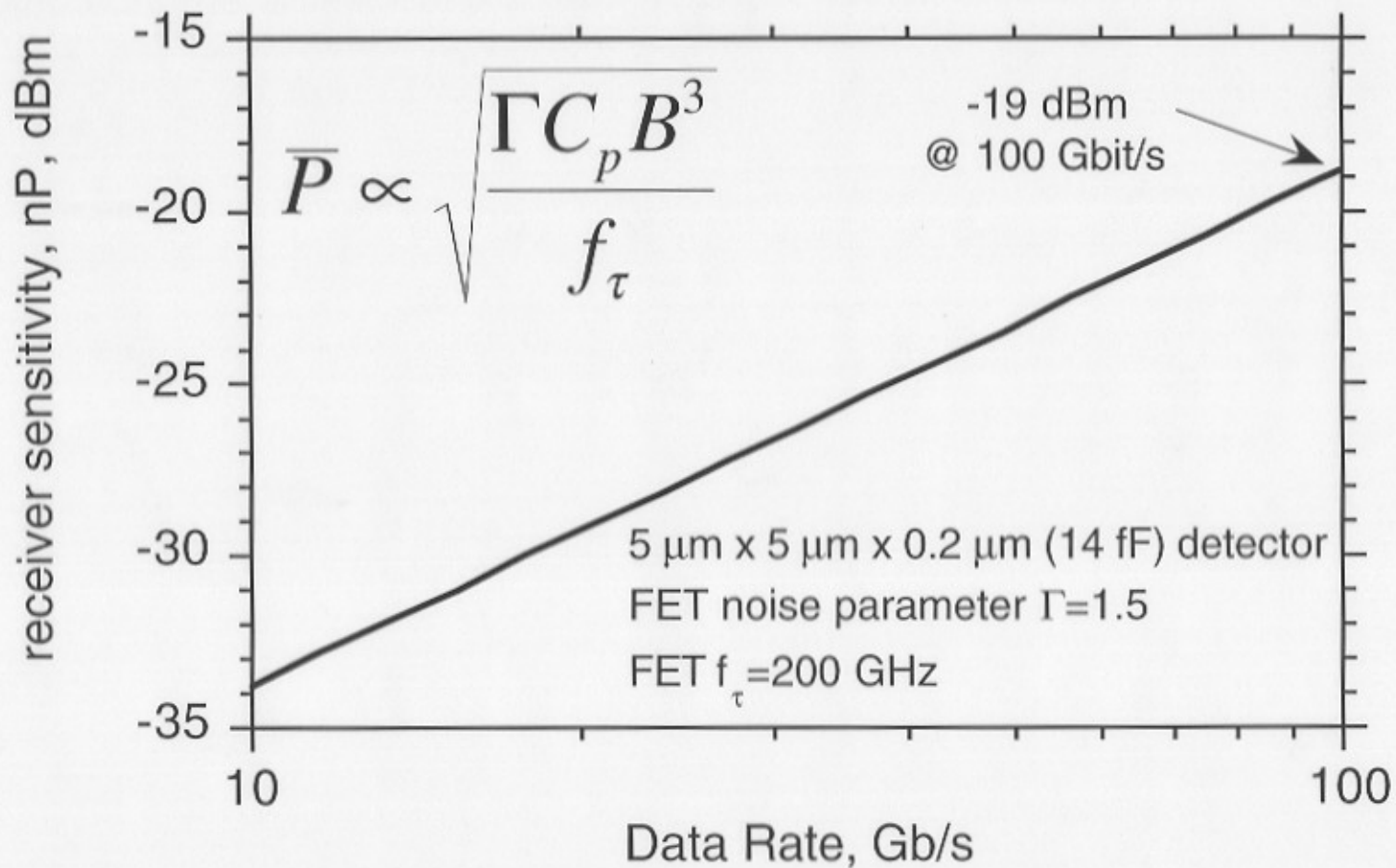
Maximize HEMT  $f_t$

Minimize detector and layout capacitance

Monolithic integration **NOT** imperative  
(well-designed bond pads add  $\approx 25$  fF capacitance)

$50\Omega$  load would incur large noise penalty

## Good transistors will make more sensitive receivers



Sensitivity can be greatly improved by using good low-noise HEMTs  
Good hybrid integration: 25 fF pad capacitance, 2-3 dB penalty.



## Calculating Sensitivity:

Channel filter *shape* with zero intersymbol interference (Nyquist filter) is assumed.

Input-referred noise current has a power spectral density of the form

$$\frac{d\langle I_n I_n^* \rangle}{df} = k_0 + k_1 f + k_2 f^2$$

Noise is amplified, and then integrated over filter bandwidth to determine RMS noise at filter output. For standard filter shapes, *Personik* has tabulated these integrals.

Dividing this by the receiver DC gain yields the input-referred RMS noise current.

Assuming Gaussian statistics, a 6:1 signal/noise ratio yields  $10^{-9}$  error rate

The average input photocurrent at sensitivity is then 6 times the frequency-integrated input-referred noise current.

## Sensitivity Expressions

Receiver sensitivity

$$\bar{P} \cong 6 \frac{h\nu}{\eta q} \bullet \sqrt{\frac{32\pi kT\Gamma C_p B^3 I_3}{f_\tau} + \frac{4kTBI_2}{R_f}}$$

B is the bit rate. For a square optical pulse and a raised-cosine eye,  $I_2=0.68$  and  $I_3=0.12$ . Sending impulses (solitons) will make  $I_3$  somewhat smaller.

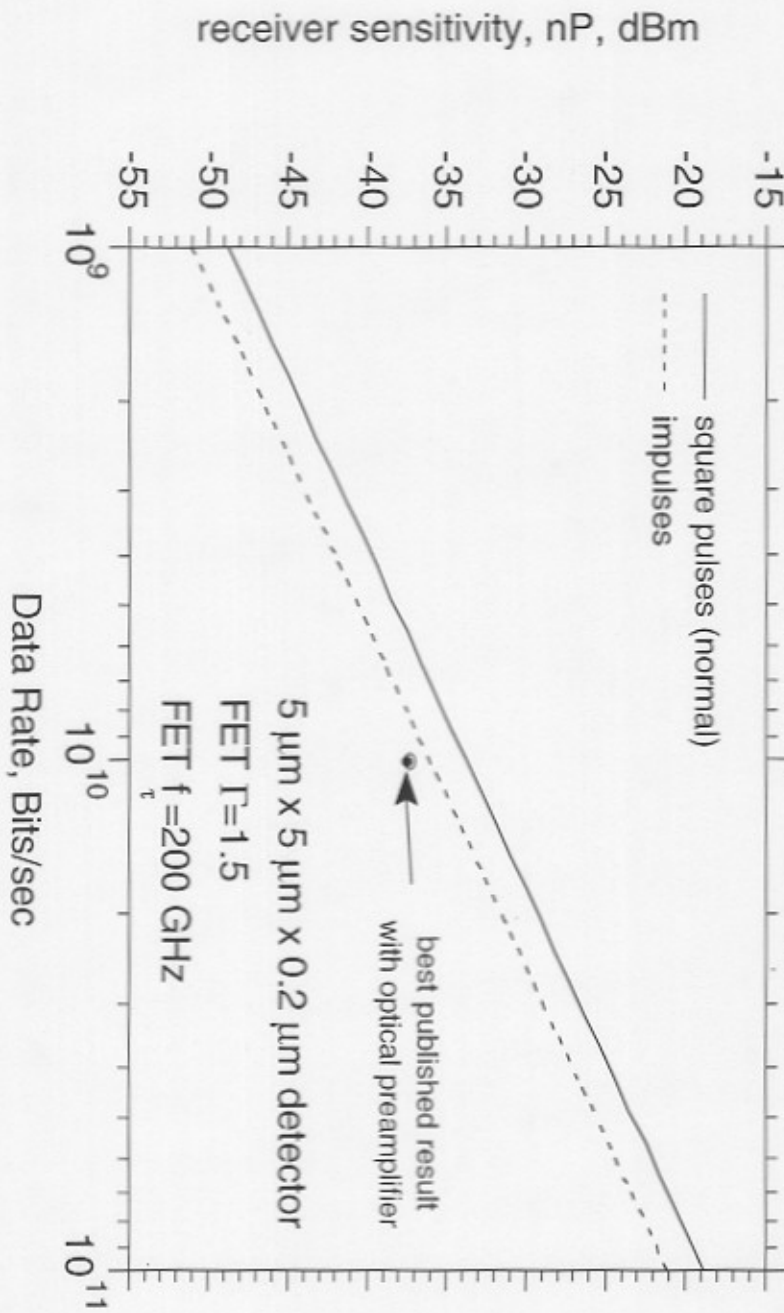
$$I_2 = [1 / BH(0)] \cdot \int H(f)df$$

$$I_3 = [1 / B^3 H(0)] \cdot \int f^2 H(f)df$$

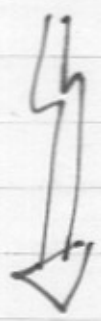
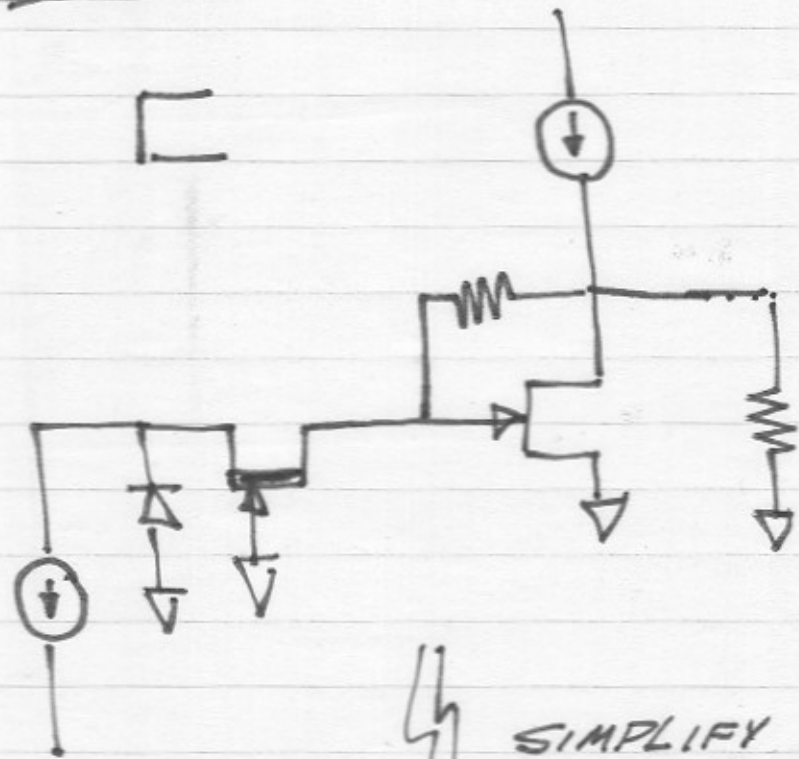
Note at 100 GB/sec, a  $5\ \mu\text{m} \times 5\ \mu\text{m} \times 0.2\ \mu\text{m}$  detector,  $f_t=200$  GHz,  $\Gamma=1.5$ , that  $R_L=209\ \Omega$  results in equal FET and resistor noise.

At 40 GB/sec, the point of equal contribution is  $R_L=1.3$  k $\Omega$ .

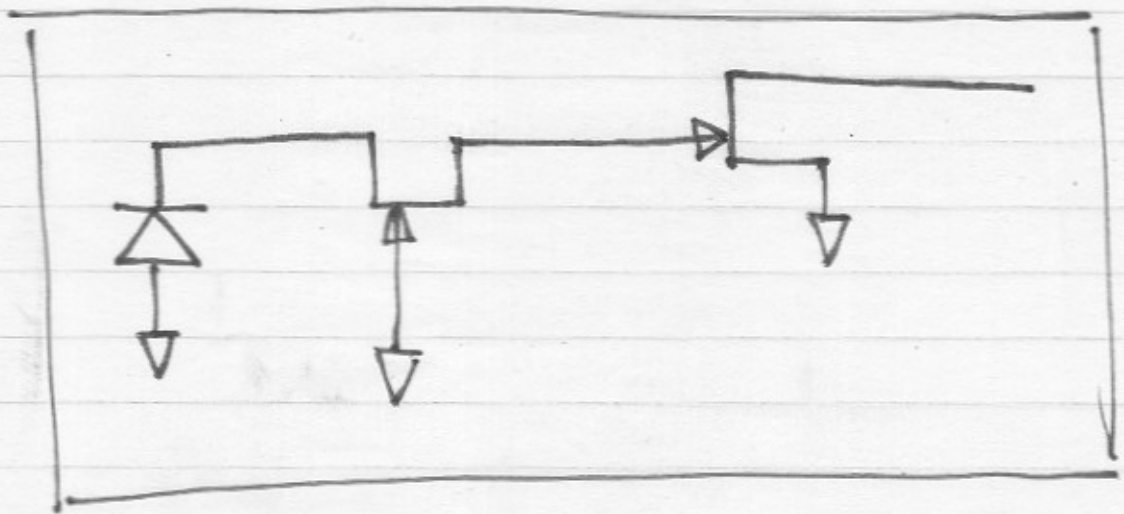
# Optical Receiver Sensitivity for square and $\delta$ -pulses



# Noise Analysis



SIMPLIFY for now  
to focus on key  
issues:



More Notes on C.G. opt. recur:

FET 2 noise, input-referred:\* ( $R_F \rightarrow \infty$ )

$$\frac{4kT\Omega}{g_{m2}} (2\pi f)^2 (C_{gs2} + g_{pd1})^2 = \frac{\sqrt{\langle II^* \rangle}}{\Delta f} \Big|_{Fet2} = A$$

$\cdot \left(1 + \left(\frac{2f}{f_T}\right)^2\right)^*$  ← ignore for now

Fet 1 noise, input-referred.

$$\frac{4kT\Omega}{g_{m1}} (2\pi f)^2 (C_{gs1} + C_p)^2 = \frac{\sqrt{\langle II^* \rangle}}{\Delta f} \Big|_{Fet1} = B$$

$$\left[ \frac{4kT\Omega}{g_{m1}} (2\pi f)^2 C_{gs1}^2 \right] 4 = \text{because } C_{gs1} = C_p \text{ minimize } B$$

Examine the ratio of second to first-stage noise, assuming  $C_{gs2} \gg C_{pd1}$

$$\frac{A}{B} = \left( \frac{C_{gs2}}{C_{gs1}} \right)^2 \frac{1}{4} \cdot \frac{g_{m1}}{g_{m2}} = \frac{1}{4} \left( \frac{W_{g2}}{W_{g1}} \right)$$

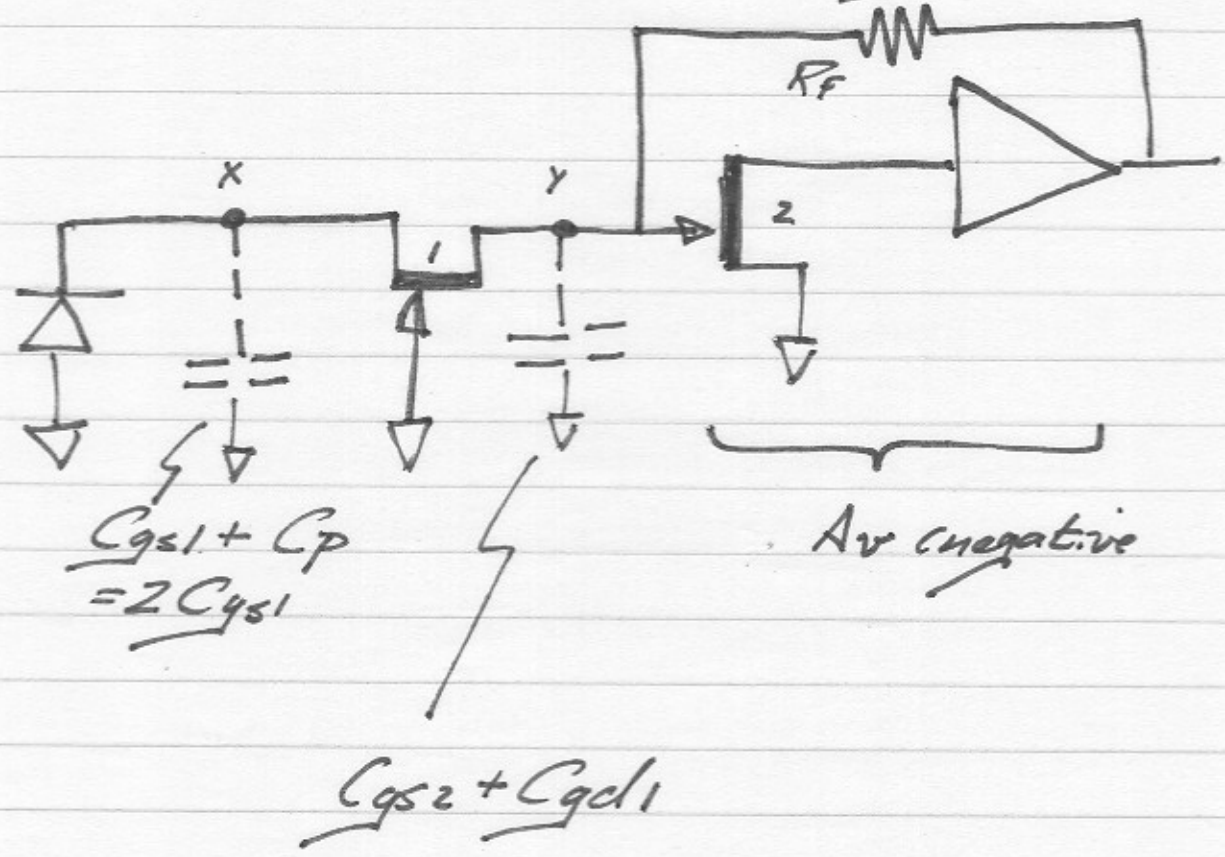
Note that  $W_{g2} \leq W_{g1}$  will ensure

negligible FET 2 noise contribution.

\* (drop for now - but check the effect later)

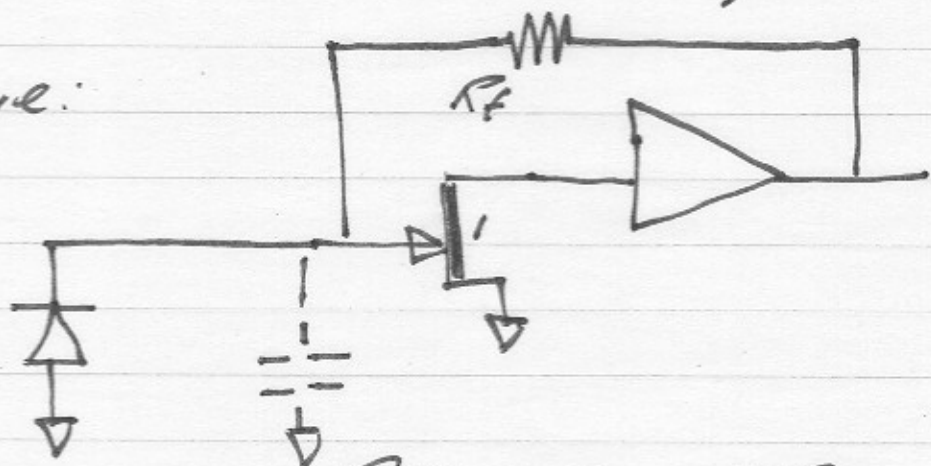
So, is that the end of the design effort?

No - we must now add the load resistor  $R_f$  and try to get adequate bandwidth without sacrificing sensitivity.



the pole at "x" is at  $f_{p1}/2$ ; while at "y" it is at  $f_{p2} = (1 - A_v) / 2\pi R_f (C_{gs2} + C_{gd1})$

Compare now to the simple transimpedance stage:



$$C_{gs1} + C_p = 2C_{gs1} = 2C_p$$

For equal values of  $R_f$ , this equal noise penalties from  $R_f$ , the trans- $Z$  has larger capacitance, (hence lower b.w.) by the ratio

$$\because C_{gd1} \ll C_{gs2}$$

$$\frac{2C_p}{C_{gs2} + C_{gd1}} = \frac{2C_{gs1}}{C_{gs2} + C_{gd1}} \approx \frac{2 \omega_{q1}}{\omega_{q2}}$$

If we make  $\omega_{q2} / \omega_{q1} = 1/2$ ,

we get four times the bandwidth with the  $C_g$  stage added.

(5)

So, it would seem that a reasonable workable solution would be to make Fet 2 have  $1/2$  the gate width of Fet #1. We will find that the noise penalty of Fet 2 is only  $1/8$  (in  $\langle I_n^2 \rangle$ , hence  $\sim 1/16$  in  $\sqrt{\langle I_n^2 \rangle}$ !) of the noise of Fet 1, and that for equal transimpedance-stage ~~to~~ bandwidths, that the resistor  $R_f$  can be four times larger.



6

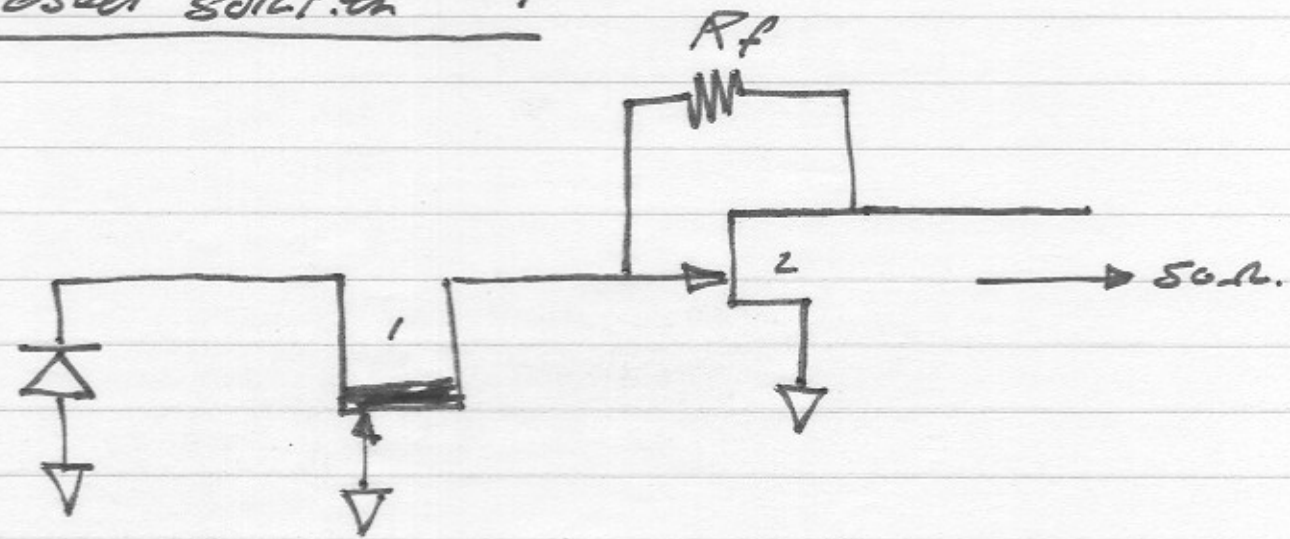
Why not simply make  $Q_2$  much smaller than  $Q_1$ , given that an "exact" analysis gives minimum  $Q_2$  noise, and bigger stage-2 bandwidth, if  $Q_2$  is very small - in fact equal to  $\omega_{q1} \times \frac{C_{q1}}{C_{q2}}$ ?

Because:

- 1) We can't make  $\omega_{q2}$  that small  
( $\omega_{q2} = \omega_{q1} \times \frac{C_{q1}}{C_{q2}}$ )
- 2) The second-stage gain becomes vanishingly small - putting a considerable burden on the noise performance of the third stage.

Instead, we should content ourselves with making the cg preamplifier four-times-better -

Proposed soln. #1



Q1:  $\frac{C_{gs1}}{W_{g1}} = C_p$

Q2:  $W_{g2} = W_{g1} / 2$

Pick  $R_f$  so that

$$\frac{4kT}{R_f} \beta I_2 = \left(\frac{1}{2}\right) \times \frac{32\pi kT \pi C_p}{f_T} \beta^3 I_3$$

the factor of 1/2 being negotiable...

8

Note that  $I_2 = 0.68$   $I_3 = 0.12$

Now a  $5\mu\text{m} \times 5\mu\text{m} \times 0.2\mu\text{m}$  detector has

$C_p = \underline{14\text{ fF}}$ , but we'll use  $25\text{ fF}$  to allow for  
stray C.

$$C_p = 25\text{ fF} \quad f_T = 100\text{ GHz}$$

$$C_{gs1} = 25\text{ fF} \rightarrow g_{m1} = 2\pi C_{gs1} f_T = 15.7\text{ mS} = 1/64\ \Omega$$

$$C_{gs2} = 12.5\text{ fF}, \quad g_{m2} = g_{m1}/2 = 1/128\ \Omega$$

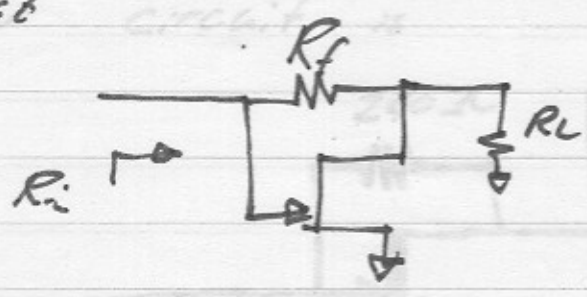
$$\frac{32\pi kT I_2 G_p B^3 I_3}{f_T} = 1.87(10^{-11})\text{ A}^2 \quad @ 100\text{ GHz.}$$

→ allowing the noise contribution of  $R_f$  to be 50% of

$$\text{this gives us } R_f = \begin{cases} 120\ \Omega & (B = 100\text{ GHz}) \\ 750\ \Omega & (B = 40\text{ GHz}) \end{cases}$$

These are minimum values.

Recall that



$$R_{in} = \frac{R_f + R_L}{1 + g_m R_L} ; \quad -Z_T = R_f \frac{g_m R_L}{1 + g_m R_L} - R_L \frac{1}{1 + g_m R_L}$$

to this point the circuit looks

100 GB/sec. Choose  $R_L = \infty$  !

& choose  $R_f = 240 \Omega$  !

$$\text{but } Z_T = R_f - \frac{1}{g_{m2}} = 240 \Omega - 128 \Omega = 112 \Omega$$

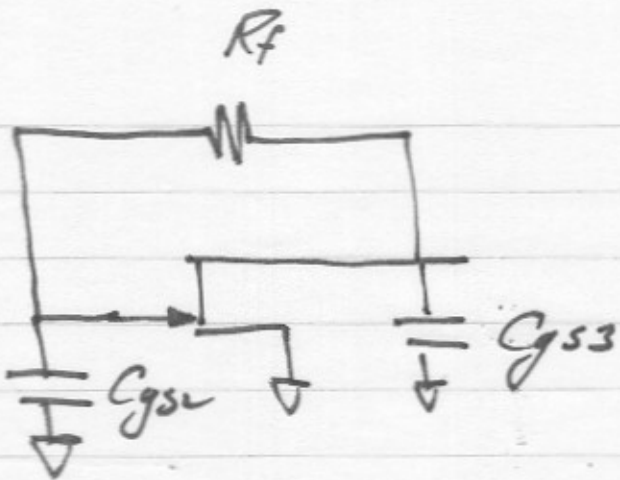
$$\text{happens } R_{in} = \frac{1}{g_{m2}} = 128 \Omega \text{ third stage poles}$$

in the second pole position: (neglecting  $C_{gd2}$ )

$$\text{for this point } f_{p2} = \frac{1}{2\pi (128 \Omega) C_{gs2}} = f_T = 100 \text{ GHz} !$$

performance gets better & better as  $R_i$  gets

larger - maybe not so.



$$a_1 = (C_{gs3} + C_{gs2}) / g_{m2}$$

$$a_2 = C_{gs3} C_{gs2} R_f / g_{m2}$$

$$\frac{a_2}{a_1} = \frac{C_{gs3} C_{gs2} R_f}{C_{gs3} + C_{gs2}} \approx C_{gs2} R_f$$

→ underdamped response will be observed

unless  $(\approx) g_{m2} R_f < 1$

→ trouble.

# Optimum Noise Design of Multi-Gigabit Common-Gate Optical Receivers

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## Abstract

### 1. Introduction

Ultra-wideband photoreceivers are needed for multi-gigabit optical transmission systems. Flat photoreceiver front-end response is desired over several decades of frequency to avoid post-equalization problems. High sensitivity at high data rates require low front-end noise at these high frequencies. The standard approach used is a front-end with a photodiode in conjunction with a transimpedance or integrating preamplifier. Several such monolithic [2] and hybrid [3] optical receivers with a p-i-n/MSM photodiode and HEMT/HBT for operation at 10-15 Gbits/sec have been reported. Excluding the use of inductive peaking [4,5] which can provide moderate noise reduction for most receiver configurations, integrating optical preamplifiers attain the lowest noise but have low bandwidth and require equalization. At low data rates, transimpedance amplifiers can attain similar noise performance. At very high data rates (>10 Gbits/sec), a transimpedance amplifier is extremely difficult to realize with the present transistor technology, as very large transistor gain-bandwidth products are demanded. Addition of a common-gate stage to the input of the transimpedance preamplifier with appropriate selection of circuit element values allows

the receiver bandwidth to be increased to one-half of the HEMT current-gain cutoff frequency, while attaining the same noise performance as the integrating front-end. Common-gate receivers were demonstrated earlier [6], but the authors presented no noise or gain-bandwidth analysis and concluded that the common-gate stage degraded the noise performance. Here, we show that, if appropriately designed, the common-gate/transimpedance based receiver can attain both wide bandwidth and very low noise. By combining this circuit configuration with advanced low-capacitance InGaAs photodetectors and AlGaAs/InGaAs/GaAs PHEMTs, optical receivers with -22 dBm sensitivity at 25 Gbits/sec and -20 dBm at 40 Gbits/sec are attainable.

## II. Analysis of the Transimpedance amplifier

In Fig. 1(a) is shown a simple transimpedance amplifier as is predominantly used in photoreceivers [7].  $R_f$  is the feedback resistor, while the input HEMT Q1, the gain stage  $A_v$  and  $R_d$  form an inverting amplifier with gain  $A = -g_{m1}R_dA_v$ . A simple first-order noise model of the HEMT (Fig. 1(b)) is generally used in optical receiver design. Neglecting all secondary noise sources and dropping non-dominant terms, the input referred equivalent noise current spectral density is given by

$$\frac{d\langle I_n I_n^* \rangle}{df} = \frac{4kT\Gamma}{g_{m1}} (2\pi f)^2 (C_p + C_{gs1})^2 + 4kT / R_f \quad (1)$$

where  $C_p$  is the photodiode capacitance,  $f_r$  is the HEMT current-gain cutoff frequency and  $\Gamma$ , its channel noise coefficient. Using the optimum HEMT size (that which gives  $C_p = C_{gs1}$ ), this becomes

$$\frac{d\langle I_n I_n^* \rangle}{df} = 32\pi kT \Gamma f^2 C_p / f_t + 4kT / R_f \quad (2)$$

At the gate of Q1, the node impedance is  $R_f / (1 - A)$  and the node capacitance is  $C_p + C_{gs1} = 2C_p$ . Hence, the preamplifier 3 dB bandwidth is given by

$$f_{3dB} = (1 - A) / 4\pi C_p R_f \quad (3)$$

We would like to make  $R_f$  large so that it contributes negligible noise. In the limit of very large  $R_f$ , we are left with the integrating front-end, which results in low front-end bandwidth and introduces significant difficulties with post-equalization. To attain large bandwidth, while still having a large  $R_f$ , the preamplifier should have a large gain  $A$ . To avoid loop instability,  $A(j\omega)$  should have no poles below the loop bandwidth  $f_{3dB}$ . Both, very high gain and wide bandwidth is difficult to attain, given finite transistor cut-off frequency  $f_t$ . Hence, multi-gigabit receivers are generally designed with  $R_f$  reduced to the point where a front-end bandwidth of about 50-75% of the data rate is attained, and sensitivity is sacrificed.

### III. Analysis of common-gate/transimpedance preamplifier

These difficulties are overcome by the common-gate/transimpedance preamplifier (Fig. 1(c)). At the source of Q1, the node impedance is  $1 / g_{m1}$  and the node capacitance is  $C_p + C_{gs1} = 2C_{gs1}$ . Hence the input pole frequency is  $f_t / 2$ . The dominant pole is at the gate of Q2, where the node impedance is  $R_f / (1 - A)$  and the node capacitance is  $C_{gd1} + C_{gs2}$ .

$$f_{3dB} = (1 - A) / 2\pi(C_{gd1} + C_{gs2})R_f \quad (4)$$



As before, taking only the dominant noise terms, the input referred equivalent noise current spectral density is given by

$$\begin{aligned} \frac{d\langle I_n I_n^* \rangle}{df} &= \frac{4kT\Gamma}{g_{m1}} (2\pi f)^2 (C_p + C_{gs1})^2 \\ &+ \frac{4kT}{R_f} (1 + (2\pi f)^2 (C_p + C_{gs1})^2 / g_{m1}^2) \\ &+ \frac{4kT\Gamma}{g_{m2}} (1/R_f^2 + (2\pi f)^2 (C_{gd1} + C_{gs2})^2) \end{aligned} \quad (5)$$

The optimum gate widths of Q1 and Q2 are those which give  $C_{gs1} = C_p$  and  $C_{gd1} = C_{gs2}$  simplifying (5) to

$$\frac{d\langle I_n I_n^* \rangle}{df} \cong 32\pi kT\Gamma f^2 C_p / f_\tau + \frac{4kT}{R_f} (1 + (2f / f_\tau)^2) + \frac{4kT\Gamma}{g_{m2} R_f^2} \quad (6)$$

where the approximation  $C_{gs} \gg C_{gd}$ , which is true for a typical HEMT, is used. From equations (2) and (6), the channel noise of HEMT Q1 is same for both configurations. If both are designed for the same bandwidth,  $R_f$  for the common-gate configuration can be much larger (by the ratio of  $C_{gs1}$  to  $C_{gd1}$ ) than for the simple transimpedance amplifier, thus giving a noise performance similar to that of an integrating front-end. For equal noise currents from both (i.e. same  $R_f$ ), the common-gate receiver will have a larger bandwidth by the ratio of  $C_{gs1}$  to  $C_{gd1}$ . In either case, the common-gate receiver works better than the simple transimpedance receiver.

#### IV. Design Illustration

We illustrate here, a design for 20-40 Gbits/sec reception (Fig. 2). Complete small-signal and noise models for the 0.3  $\mu\text{m}$  gate length pseudomorphic HEMTs are taken from [8] while a 100 GHz 5 $\mu\text{m}$  x 5 $\mu\text{m}$  GaInAs/InP p-i-n photodiode [9] is assumed. The widths of

the HEMTs Q1 and Q2 are chosen for low input noise. The feedback resistors are chosen for large bandwidth while not degrading the noise performance. Q1 is the common-gate stage while Q2 and  $R_f$  form the transimpedance stage. Q3 and Q4 provide both, voltage gain and buffering to drive a  $50 \Omega$  load.

Fig. 3 shows SPICE simulations of transimpedance gain and input noise current using complete device models. The simulation includes 25 fF parasitic assembly capacitance consistent with MMIC assembly. The 3 dB bandwidth is 20 GHz and the input referred noise current is 14 pA/sqrt. Hz at 20 GHz. The predicted receiver sensitivity [7] at  $10^{-9}$  BER (Fig. 4) is -19 dBm at 40 Gbits/sec. For comparison, the predicted sensitivity of an integrating front-end is also shown. It can be seen that the common-gate receiver approximates the minimum attainable noise of an integrating front-end.

## V. Conclusions

We have analyzed the simple transimpedance preamplifier and the common-gate/transimpedance preamplifier for use in optical receiver front-ends. Our analysis shows that, for the same bandwidth of the two circuits, the common-gate-transimpedance preamplifier is less noisy and hence, more sensitive. On the other hand, if both are designed for the same noise performance, the common-gate/transimpedance preamplifier has a higher gain-bandwidth. In the case of the common-gate/transimpedance preamplifier, the trade-off between noise and bandwidth is less than in the case of the simple transimpedance amplifier. SPICE simulations of our front-end design using state-of-the-art devices predict satisfactory performance for 20-40 Gbits/sec. In conclusion, we would like to say that the common-gate/transimpedance preamplifier is a promising candidate for high bitrate photoreceivers.

## References

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- [8] Michael Schlechtweg, Werner Reinert, Paul J. Tasker, Ronald Bosch, Jurgen Braunstein, Axel Hulsmann and Klaus Kohler, "Design and Characterisation of High Performance 60 GHz Pseudomorphic MODFET LNAs in CPW-Technology Based on

Accurate S-Parameter and Noise Models," IEEE Trans. on Microwave Theory and Techniques, vol. 40, no. 12, p. 2445, 1992

[9] Yih-Guei Wey, Kirk S. Giboney, John E. Bowers and Mark J. W. Rodwell, "100 GHz Double Heterostructure GaInAs/InP p-i-n Photodiode," OSA Proceedings on Ultrafast Elect. and Optoelect., vol. 14, 1993

## Figure Captions

Fig. 1 Simple transimpedance amplifier (a), Simplified HEMT noise model (b), and Common-gate/transimpedance preamplifier (c)

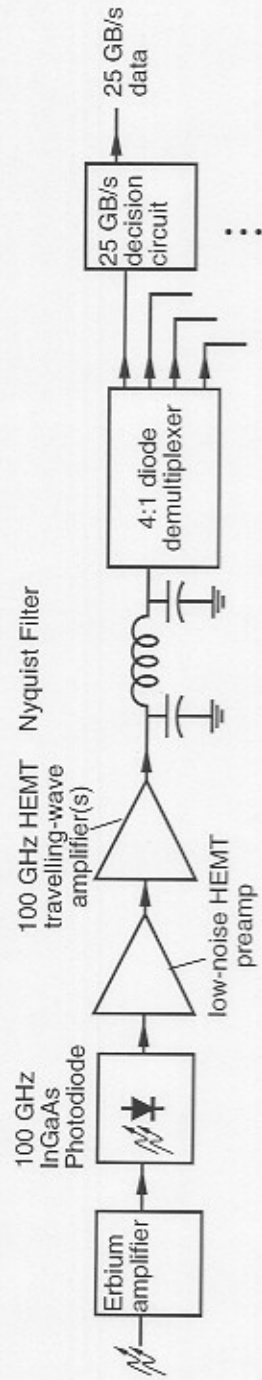
Fig. 2 Schematic diagram of common-gate/transimpedance photoreceiver

Fig. 3 SPICE simulations of transimpedance gain and input referred equivalent noise current spectral density

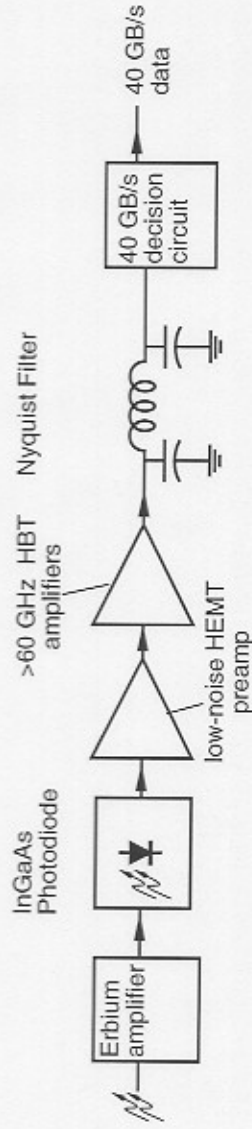
Fig. 4 Predicted sensitivity of common-gate/transimpedance preamplifier and integrating front-end using Personick coefficients  $I_2=0.68$  and  $I_3=0.12$

# Receiver Implementation: In Summary

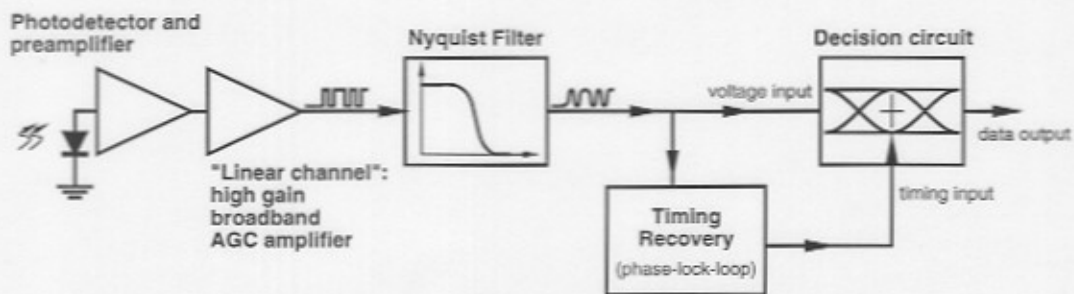
## 100 GB/sec Optical Receiver

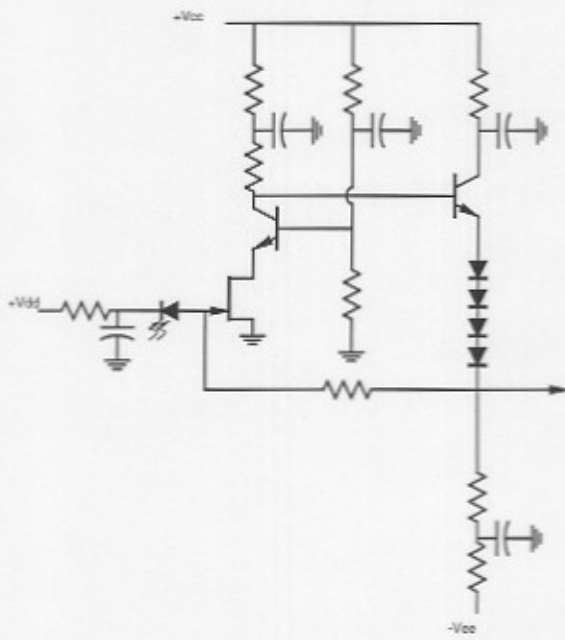


## 40 GB/sec Optical Receiver



*Block diagram*





Front End,



Sensitivity Comparison, HBT phototransistor vs HBT-PIN, PIN-FET and APD-FET

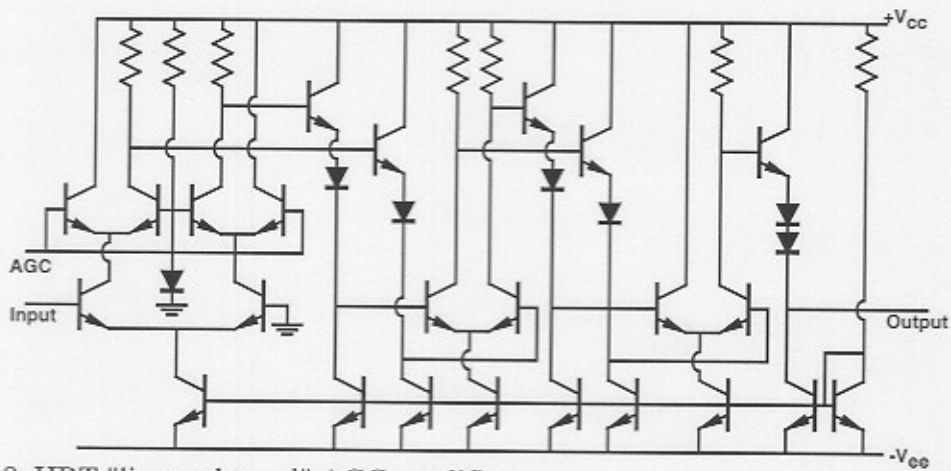
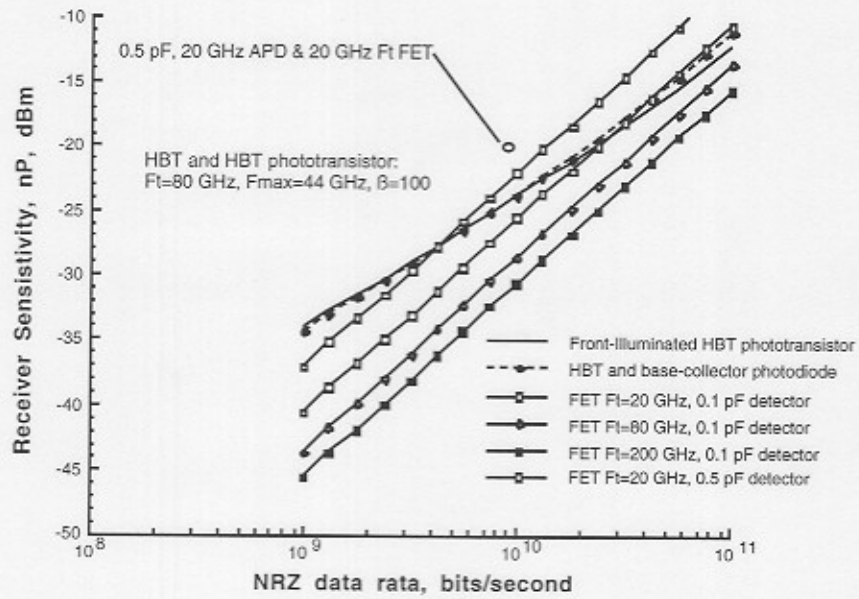
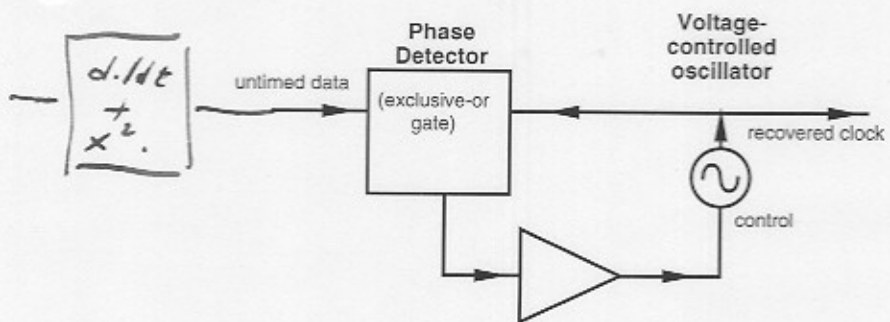


Figure 9. HBT "linear-channel" AGC amplifier

# Phase lock loop.



Decision Circuit

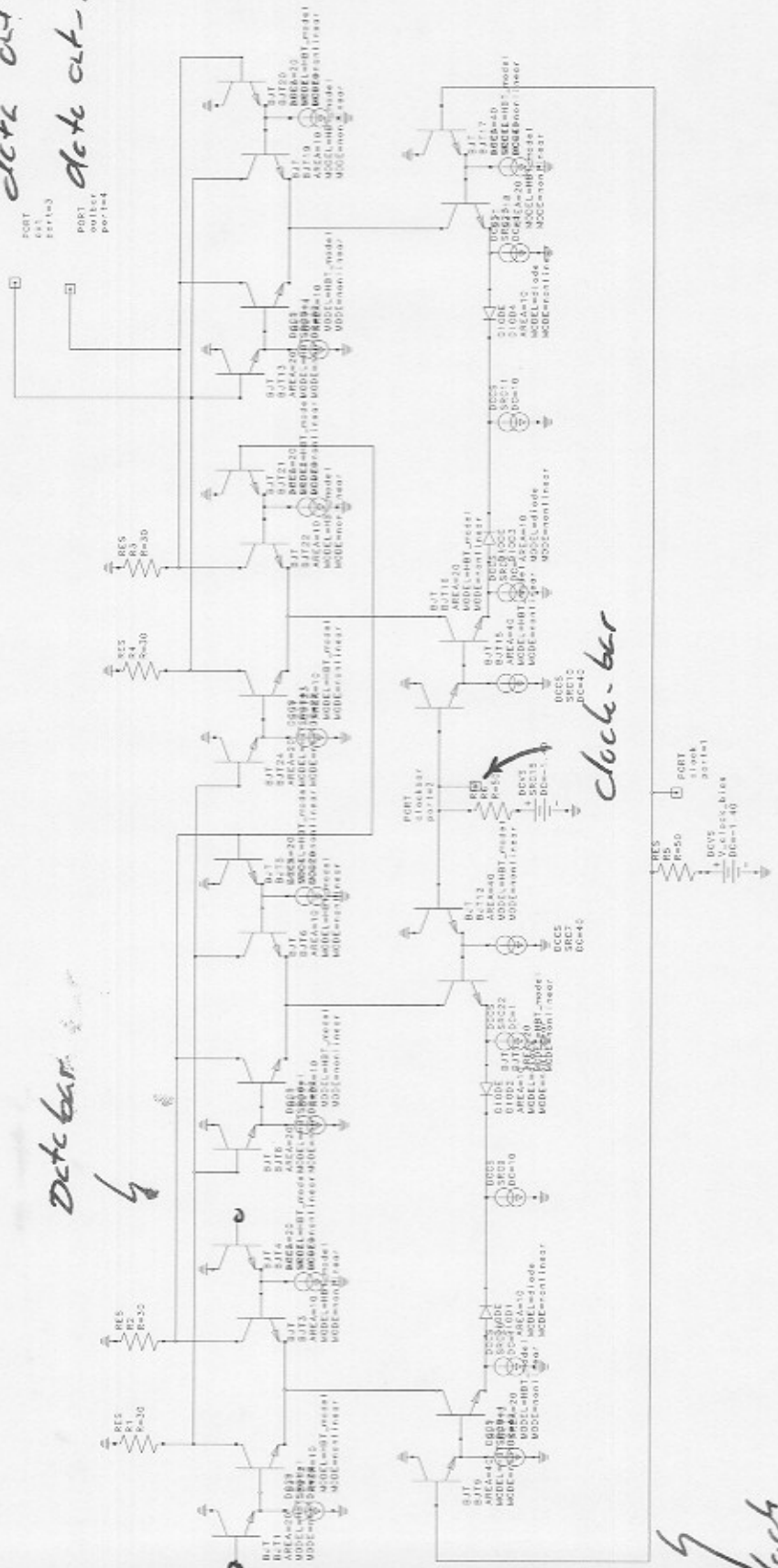
data out  
data out-bar

data bar

data

clock-bar

clock



# AGC Amplifier

gain control

outputs

acg detectors

