

ECE 162A

Mat 162A

Lecture #2: Wavelike Properties of
Matter

Read Chapter 3 of Eisberg, Resnick

- What is the classical theory of light?

Maxwell's Equations

$$\nabla \times H = I + \frac{\partial D}{\partial t}$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\nabla \cdot D = 0$$

$$\nabla \cdot B = 0$$

where E and H are the electric and magnetic field vectors
 D and B are the electric and magnetic displacement vectors
No free charge. No current flow.

Constitutive Relations

$$D = \epsilon_0 E + P$$

$$B = \mu_0 (H + M)$$

P and M are the electric and magnetic polarizations of the medium
 ϵ_0 and μ_0 are the electric and magnetic permeabilities of vacuum
 E and H are the electric and magnetic field vectors
 D and B are the electric and magnetic displacement vectors

For isotropic media, Electric Susceptibility χ

Isotropic Media: χ is a complex number

$$P = \epsilon_0 \chi E$$

The real part determines the index (velocity) and the imaginary part determines the gain or absorption.

Isotropic media: Vacuum, gasses, glasses (optical fibers)

Anisotropic media: Semiconductors, crystalline materials.

Wave Propagation in Lossless, Isotropic Media

- Lossless: $\sigma=0$, χ is real, ϵ is real.
- Isotropic: χ , ϵ are scalars (not tensors).

$$\nabla \times E = i + \frac{\partial B}{\partial t} = 0 + \mu \frac{\partial H}{\partial t}$$

$$\nabla \times H = i + \frac{\partial D}{\partial t}$$

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$$\nabla \times H = i + \frac{\partial D}{\partial t}$$

$$\nabla \times (\nabla \times E) = \mu \frac{\partial(\nabla \times H)}{\partial t} = \mu \frac{\partial^2 D}{\partial^2 t} = \mu\epsilon \frac{\partial^2 E}{\partial^2 t}$$

$$\nabla \times (\nabla \times E) = \nabla^2 E - \nabla(\nabla \cdot \vec{e})$$

Wave Equation

$$\nabla^2 E = \mu\epsilon \frac{\partial^2 E}{\partial^2 t}$$

Wave Equation

$$E(x, y, z, t) = \text{Re}[E(x, y, z)e^{i\omega t}]$$

$$\nabla^2 \vec{E} + \omega^2 \mu \epsilon \vec{E} = 0$$

$$\nabla^2 \vec{E} + k^2 \vec{E} = 0$$

where

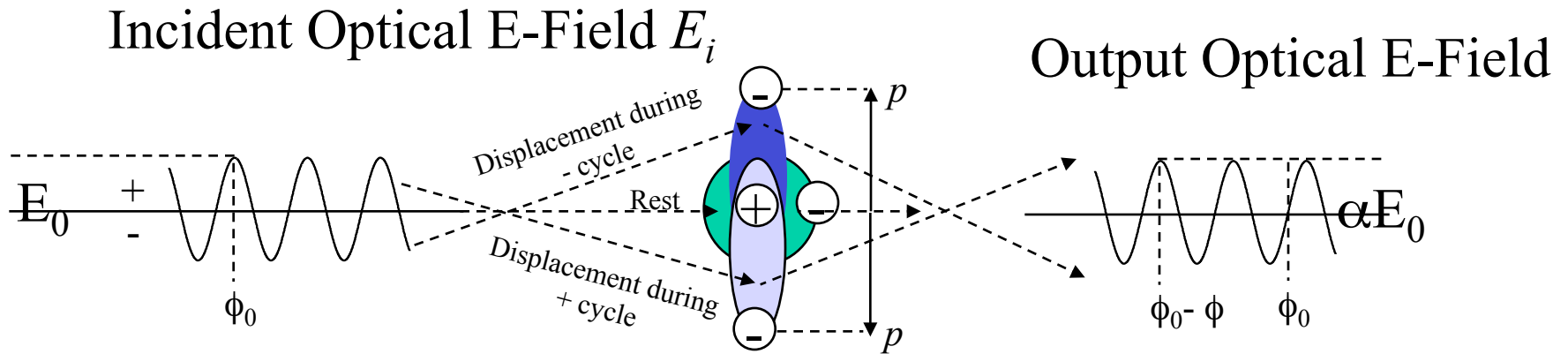
$$k = \omega \sqrt{\mu \epsilon} = \omega n / c \quad \text{Wave number}$$

$$c = 1 / \sqrt{\mu_0 \epsilon_0} \quad \text{Speed of light}$$

$$n = \sqrt{\frac{\mu \epsilon}{\mu_0 \epsilon_0}} \quad \text{Index of refraction}$$

Radiation and Atomic Systems

In a dielectric material, an incident optical E-field oscillating at frequency ω will induce physical displacement of a bound electron from its rest state at the same oscillation frequency. We call the the *induced dipole moment or displacement*



α is field attenuation

p is the induced dipole moment

α_p is the atomic polarizability

ϕ_0 is the input field phase reference

ϕ is the output field phase

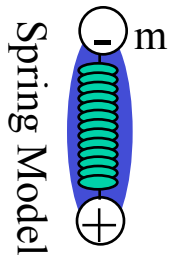
$$\mathbf{p} = \alpha_p \mathbf{E}_i$$

Classical Electron Model

- Now we need a basic model that describes the motion, or displacement of the electron (dipole) in the presence of an external force (incident optical field)
- Consider the classical equation of motion that describes displacement (position X) of an electron with charge q and mass m attached to an atom (as if connected by a spring) with a dampening coefficient (loss) γ and resonant frequency ω_0 in response to an applied field

$$m \frac{d^2}{dt^2} X + m\gamma \frac{d}{dt} X + m\omega_0^2 X = -qE_a$$

- ⇒ Lets assume the atom can be described as a two level system with energy levels E_1 and E_2 such that the atomic resonance is $\omega_0 = (E_2 - E_1)/\hbar$ and the applied electric field is of the form $E = E_0 e^{i\omega t}$. Then the electron position (relative to rest) and induced dipole moment of the atom can be described as



$$X = \frac{-qE_0}{m(\omega_0^2 - \omega^2 + i\gamma\omega)} e^{i\omega t}$$

$$p = -qX = \frac{-q^2 E_0}{m(\omega_0^2 - \omega^2 + i\gamma\omega)} e^{i\omega t}$$

Classical Electron Model

- The dielectric constant can now be written in terms of the spring model

$$p = \alpha_p E_a$$

$$\alpha_p = \frac{q^2}{m(\omega_0^2 - \omega^2 + i\gamma\omega)}$$

$$n^2 = 1 + \chi = 1 + \frac{N\alpha_p}{\epsilon_0} = 1 + \frac{Nq^2}{m\epsilon_0(\omega_0^2 - \omega^2 + i\gamma\omega)}$$

- ⇒ And we can now re-write the electric susceptibility in terms of the classical electron model and define the real and imaginary parts of χ as χ' and χ''

$$\chi = \chi' - i\chi'' = \frac{Nq^2}{m\epsilon_0(\omega_0^2 - \omega^2 + i\gamma\omega)}$$

Near Resonance Condition

- Near resonance, $\omega \approx \omega_0$, the second term in χ is small compared to the first term and the index of refraction can be approximated by $n=1+ \chi/2$

$$\begin{aligned}n &\approx 1 + \frac{\chi}{2} = 1 + \frac{Nq^2}{2m\epsilon_0(\omega_0^2 - \omega^2 + i\gamma\omega)} \\ &= 1 + \frac{Nq^2}{4\omega_0 m\epsilon_0(\omega_0 - \omega + i\gamma/2)}\end{aligned}$$

Complex Refractive Index and Dispersion

- The real and imaginary parts of the electric susceptibility and index of refraction give rise to *optical phase delay* and *optical absorption*.
- ⇒ We can separate the real and imaginary parts by multiplying numerator and denominator by the complex conjugate

$$\begin{aligned}
 \chi' - i\chi'' &= \frac{Nq^2}{m\varepsilon_0(\omega_0^2 - \omega^2 + i\gamma\omega)} \frac{\omega_0^2 - \omega^2 - i\gamma\omega}{\omega_0^2 - \omega^2 - i\gamma\omega} \\
 &= \frac{Nq^2(\omega_0 - \omega)}{2m\omega_0\varepsilon_0[(\omega_0 - \omega)^2 + (\gamma/2)^2]} - i \frac{Nq^2(\gamma/2)}{2m\omega_0\varepsilon_0[(\omega_0 - \omega)^2 + (\gamma/2)^2]} \\
 \chi' &= \frac{Nq^2(\omega_0 - \omega)}{2m\omega_0\varepsilon_0[(\omega_0 - \omega)^2 + (\gamma/2)^2]} \\
 \chi'' &= \frac{Nq^2(\gamma/2)}{2m\omega_0\varepsilon_0[(\omega_0 - \omega)^2 + (\gamma/2)^2]}
 \end{aligned}$$

Complex Refractive Index and Dispersion

- The complex index is (not assuming near resonance)

$$n' = n - i\kappa = 1 + \frac{Nq^2(\omega_0^2 - \omega^2)}{2m\epsilon_0 [(\omega_0^2 - \omega^2)^2 + (\gamma^2\omega^2)]} - i \frac{Nq^2\gamma\omega}{2m\epsilon_0 [(\omega_0^2 - \omega^2)^2 + (\gamma^2\omega^2)]}$$

- ⇒ For a propagating field, we can write

$$\begin{aligned} E &= Ae^{i(\omega t - k'z)} = Ae^{i(\omega t - \frac{2\pi}{\lambda}(n - i\kappa)z)} \\ &= Ae^{i(\omega t - \frac{2\pi}{\lambda}nz)} e^{-\frac{2\pi}{\lambda}\kappa z} \end{aligned}$$

- ⇒ Where we can now define the attenuation coefficient α

$$\alpha = \frac{1}{I} \frac{dI}{dz}$$

$$I(z) = I(0)e^{-\alpha z}$$

$$\alpha = \frac{4\pi}{\lambda}\kappa$$

Dispersion and Complex Refractive Index

- Looking more closely at n from previous lecture (harmonic oscillator model),

$$n \approx 1 + \frac{\chi}{2} = 1 + \frac{Nq^2}{2m\epsilon_0(\omega_0^2 - \omega^2 + i\gamma\omega)}$$

$$= 1 + \frac{Nq^2}{4\omega_0 m\epsilon_0(\omega_0 - \omega + i\gamma/2)}$$

- Using the notation for n as n' to denote a complex number, we see that as ω approaches ω_0 and assuming $\chi \ll 1$, the index of refraction increases. This change in refractive index as a function of optical frequency is called **Chromatic Dispersion**.
- Looking at the imaginary part of n' we see that the optical E-field will be attenuated as ω approaches ω_0 giving rise to **Optical Absorption**.

$$n' = n - i\kappa = 1 + \frac{Nq^2(\omega_0^2 - \omega^2)}{2m\epsilon_0[(\omega_0^2 - \omega^2)^2 + (\gamma^2\omega^2)]} + i \frac{Nq^2\gamma\omega}{2m\epsilon_0[(\omega_0^2 - \omega^2)^2 + (\gamma^2\omega^2)]}$$

Chromatic Dispersion
Absorption

Dielectric Constant and Index of Refraction

- For a homogeneous material of N atoms per unit volume, the polarization is defined by the contribution of all atoms in that volume that interact with the optical E-field

$$\mathbf{P} \approx N\mathbf{p} = N\alpha_p \mathbf{E}_i \equiv \varepsilon_0 \chi_e \mathbf{E}_i$$

⇒ Where we have defined the *vacuum permittivity* ε_0 and *material electric susceptibility* χ_e by

$$N\alpha_p \equiv \varepsilon_0 \chi_e$$

⇒ Using the constitutive relation $\mathbf{D} = \varepsilon \mathbf{E} = \varepsilon_0 \mathbf{E} + \mathbf{P}$, the *dielectric constant* is defined as

$$\begin{aligned} \varepsilon &= \varepsilon_0 + \frac{\mathbf{P}}{\mathbf{E}} = \varepsilon_0 + \varepsilon_0 \chi_e = \varepsilon_0 (1 + \chi_e) \\ &= \varepsilon_0 \left(1 + \frac{N\alpha_p}{\varepsilon_0} \right) \end{aligned}$$

⇒ And we can now define for a non-magnetic medium (which dielectrics are at optical frequencies), the *optical index of refraction* n

$$\begin{aligned} n^2 &= 1 + \chi_e = 1 + \frac{N\alpha_p}{\varepsilon_0} \\ n &= \sqrt{1 + \chi_e} = \sqrt{1 + \frac{N\alpha_p}{\varepsilon_0}} \end{aligned}$$

Prequantum Theory

Chapter 1: Thermal Radiation, Plank's Constant

- Light is a wave. Maxwell's equations give rise to a wave equation that explain light propagation quite well. (Undergrad E&M)
- Classical:
 - Wavelength may be quantized (satisfying boundary conditions)
 - Wave can have any energy (continuous)
- Classical theory predicts diffraction, refraction, propagation very well.
- Electromagnetic radiation spreads through space like water waves spread across water.

- What is wrong with classical theory of light?

Chapter 2: Light: particle and wave characteristics

- Photons: Particle like properties of radiation
- Interaction of light with matter:
 - Photoelectric effect
 - Compton effect
 - Pair production
 - Bremsstrahlung
 - Pair annihilation
- All show experimental evidence of particle nature of light when interacting with matter.

What are these effects?

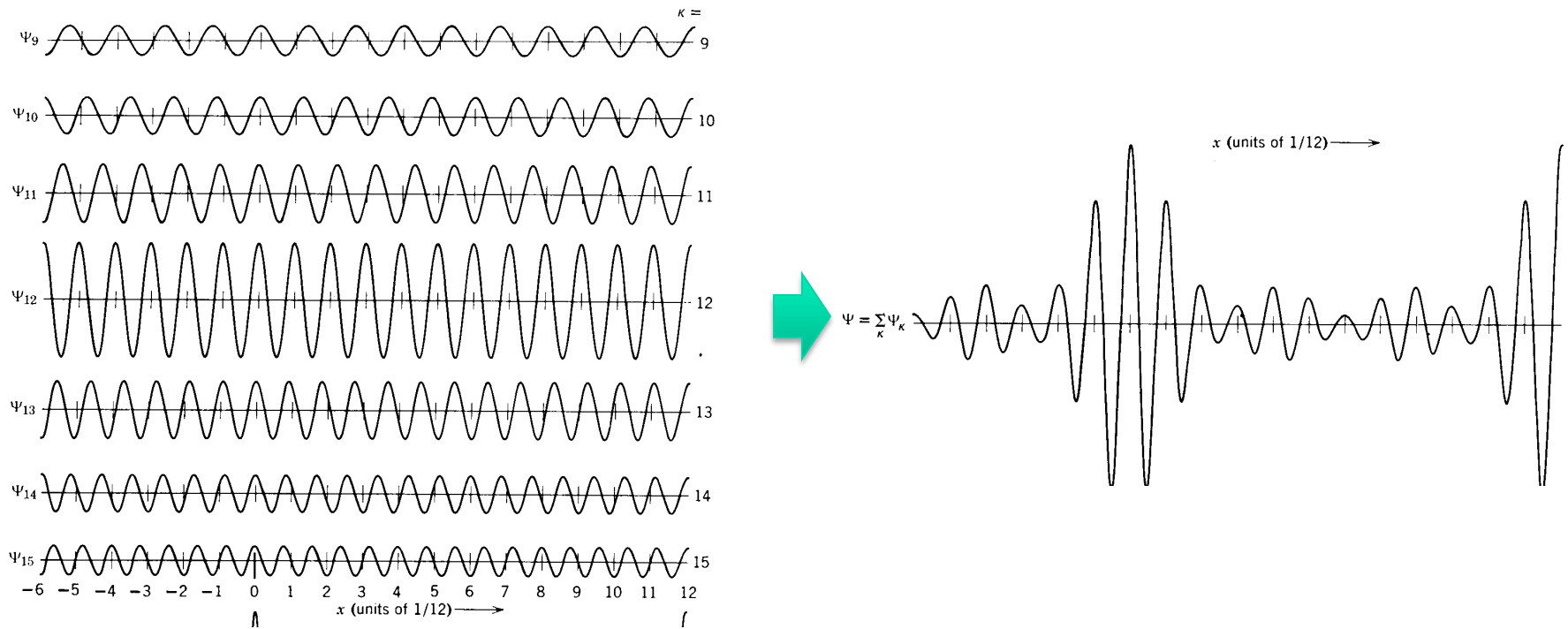
Dual Nature of Radiation

- **Radiation is neither purely a wave phenomenon nor purely a particle phenomenon.**
- A crystal spectrometer used to measure X ray wavelength is using the wave nature.
- A Compton scattering experiment is characterizing the particle nature.

- OK. So, light has both wave and particle nature.
- What is a photon?

So, what is a photon?

- A superposition of cw waves added in phase such that the envelope is a pulse.



What does a 2 slit diffraction pattern look like?

What happens when you send a single photon through a slit?

What happens if you then send 1 million photons one at a time through the slit?

Chapter 3: Wave Nature of Matter

- de Broglie: Ph.D. thesis: Wave nature of matter.
- The dual wave-particle behavior of light applies equally well to matter.
- Received the Nobel prize 5 years later.
- Energy: $E=h\nu$
- Momentum: $p=h/\lambda$
- Wavelength: $\lambda=h/p$

Examples

- What is the wavelength of a fastball?
 $m=1 \text{ kg. } v=10 \text{ m/s}$
- What is the wavelength of a 100 eV electron?
(used in MBE for low energy electron scattering)

$$h=6.6 \times 10^{-34} \text{Js}$$

Matter Wavelengths

- Baseball wavelength:
 - $\lambda = h/p = h/(mv) = 6.6 \times 10^{-34} / (1 \text{ kg } 10 \text{ m/s}) = 6.6 \times 10^{-35} \text{ m} = 10^{-25} \text{ angstrom}$
- Electron wavelength: ($E = p^2/2m$)
 - $5.4 \times 10^{-24} \text{ kg m/s}$
 - $\lambda = h/p = 6.6 \times 10^{-34} / 5.4 \times 10^{-24} \text{ m/s} = 10^{-10} \text{ m} = 1 \text{ \AA}$

Validation of matter waves: Davisson/Germer experiment

- Electrons incident on a single crystal sample (nickel).
- Angular dependence only explainable through diffraction not particle scattering

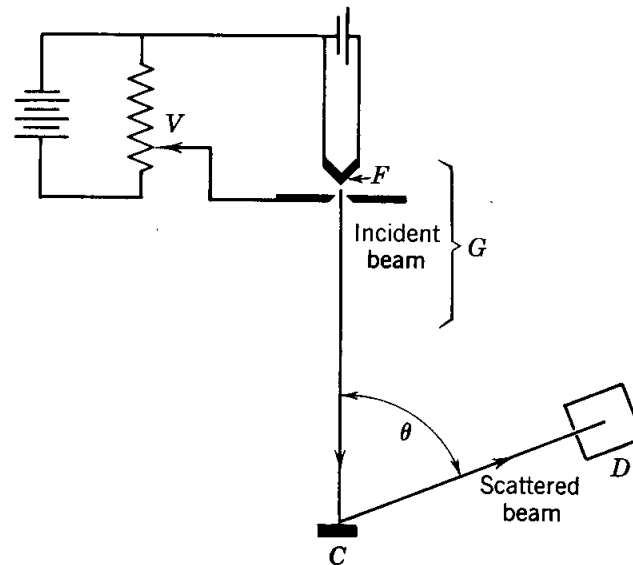


Figure 3-1 The apparatus of Davisson and Germer. Electrons from filament F are accelerated by a variable potential difference V . After scattering from crystal C they are collected by detector D .

X Ray Diffraction Electron Diffraction

- Bragg relation:
- $n\lambda = 2d \sin \varphi$

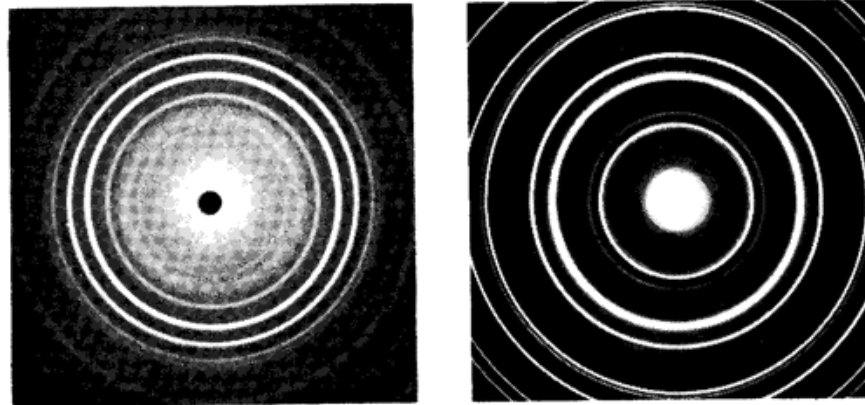
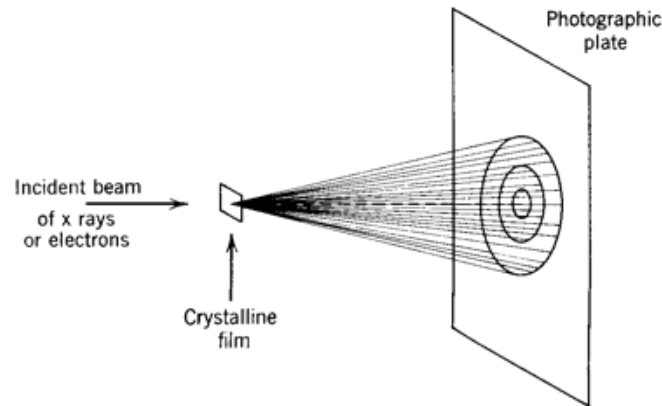


Figure 3-4 Top: The experimental arrangement for Debye-Scherrer diffraction of x rays or electrons by a polycrystalline material. Bottom left: Debye-Scherrer pattern of x-ray diffraction by zirconium oxide crystals. Bottom right: Debye-Scherrer pattern of electron diffraction by gold crystals.

Classical particle motion

- For large objects, the mass is large so the wavelength is very small, smaller than can be observed experimentally.
- $\lambda = h/p = h/(mv)$

Principle of Complementarity

- Neils Bohr: The wave and particle models are complementary; if a measurement proves the wave character of radiation or matter, then it is impossible to prove the particle character in the same experiment.
- Which model is used (wave or particle) is determined by the experiment.

Light Particle/Wave Link

- Link provided by a probability interpretation.
- Wave: Intensity proportional to average electric field squared.
- Particle: Intensity equal to $Nh\nu$ where N is the average number of photons per unit time per unit area (perpendicular to propagation vector).
- $I = \overline{E^2}/(\mu_0 c) = N h \nu$

Matter Particle/Wave Link

- Link provided by a probability interpretation.
- Particle: Probability density is proportional to average wave function squared.

Classical and Quantum Physics

- Classical physics: Laws of physics are deterministic. The laws of motion can be solved exactly and the position and momentum can be known for all time.
- Quantum physics: Laws of physics determine probability of finding a particle. The position and momentum cannot be known exactly at any time, much less all time.
- (Einstein famous objection: God does not play dice with the universe).

Fourier Analysis

- A gaussian transforms to a gaussian.
- Spectral analysis: A gaussian pulse in time is composed of a variety of frequencies, with an envelope that is a gaussian.
- Narrow frequency distribution means large time distribution
- Narrow pulse in time requires a large range of frequencies.

$$\Delta t \Delta \omega > \frac{1}{2}$$

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$$\Delta t \Delta \omega > \frac{1}{2}$$

If $E = h\nu = h\omega$, then

$$\Delta t \Delta E > \frac{1}{2} h$$

Heisenberg Uncertainty Principle

- One cannot simultaneously measure energy and time better than

$$\Delta t \Delta E > \frac{1}{2} \hbar$$

- One cannot simultaneously measure moment and position better than

$$\Delta p_x \Delta x > \frac{1}{2} \hbar$$

$$\Delta p_y \Delta y > \frac{1}{2} \hbar$$

$$\Delta p_z \Delta z > \frac{1}{2} \hbar$$

We will define Δx more carefully later. For now, it is the uncertainty in position.

- An attempt to measure position accurately requires high energy light (for example) which makes the momentum uncertain.

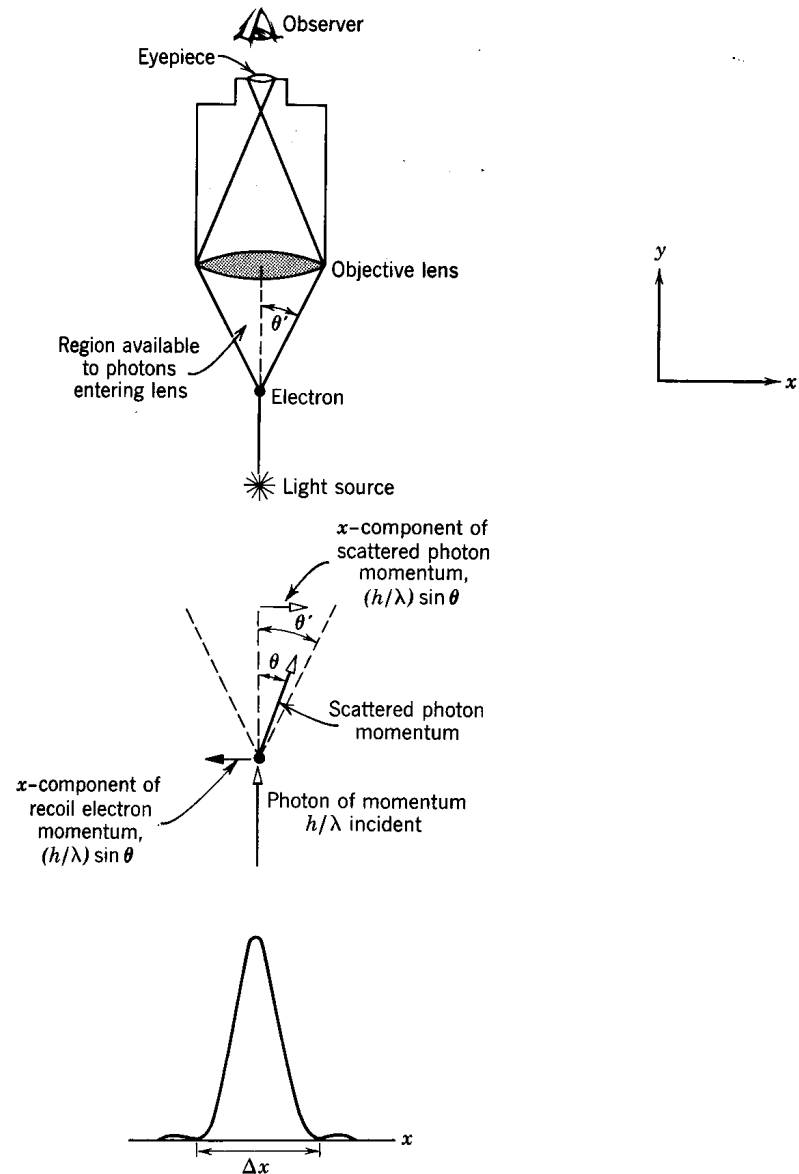


Figure 3-6 Bohr's microscope thought experiment. *Top:* The apparatus. *Middle:* The scattering of an illuminating photon by the electron. *Bottom:* The diffraction pattern image of the electron seen by the observer.

Particle/Wave Duality

- All material objects show both particle and wave aspects.
- The uncertainty principle means that an experiment to determine particle aspects (for example position) means that momentum is unknown (i.e. wavelength is unknown) and vice versa.

- Assignment #1: Assigned Tuesday Oct. 6th on the web
- No class next Tuesday
- Recitation only on Thursday Oct. 8th
- Makeup class Friday Oct. 9th
- Assignment #1 due Tuesday 13th