

#### Lecture 3: Photodetectors

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### Photodetectors (Continued)

#### $\Rightarrow$ Last lecture we covered

- photoconductors and the dominant photon absorption mechanism: Intrinsic (Band-to-Band)
- ⇒ the power absorbed as a function of wavelength for different materials and derived the efficiency  $\eta(\lambda, x)$
- ⇒ the concept of carrier lifetime and transit time and the resulting photoconductive gain (G) that results from a mismatch in electron and hole mobilities
- ⇒ Frequency response of photoconductive photodetectors

### Biased p-n Junction Photodiodes



### p-n Junction Photodiode Equation



•  $I_{dark}$  = is the current that occurs with zero optical input

 $\bullet I_s = I_{th}$  is the thermal or saturation current that occurs in normal (non-illuminated) diode operating mode

- $I_{\text{photo}}$  is photo-generated current =  $\frac{\eta q}{hv} P_{rcvd}$
- q is the electron charge
- $V_{bias}$  is applied bias voltage (positive = forward, negative=reverse)
- K<sub>B</sub> is Boltzman's constant
- T is temperature (usually in Kelvin, depending on units of  $K_B$ )

## p-n Junction Photodiode Regions of Operation



### pn-Junction Carrier Dynamics (1)

- ⇒ Carrier diffusion time (~ns/µm) is typically much longer than carrier transit time (~10ps/µm)
- ⇒ Electron and hole velocities saturated in depletion region due to high field strength
- ⇒ Once away from depletion region carrier velocities fall below saturation
- ⇒ Space charge barrier prevents carriers from entering the depletion region, therefore the multiple carrier effect seen in photoconductors does not occur when carrier velocities are mismatched



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Lecture 3, Slide 6

### pn-Junction Carrier Dynamics (2)

- Photons absorbed within one diffusion length outside the depletion region will be absorbed and the current contributing carriers will suffer both diffusion time and transit time delays
- ⇒ Effect is geometry and material dependent



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Lecture 3, Slide 7

### pn-Junction Carrier Dynamics (3)

⇒ The separation of charge in the depletion region (due to uncompensated Donors and Acceptors) leads to a capacitive effect that also impacts the detector bandwidth





⇒ The frequency at which the detector bandwidth rolls off by 3-dB due to the junction capacitance is

$$\omega_{RC} \frac{1}{R_S C_j}$$

Cj = area of depletion region  $l_d$  = depletion region length

### p-i-n Photodiodes

- To increase the photon absorption region, a layer of *intrinsic* semiconductor material can be added beteween the p and n material.
- The pin photodetector gain-bandwidth product improves of the pn-junction
  - The detector quantum efficiency can be increased over that of a simple pn junction since the depletion region is almost entirely contained in the intrinsic region and the intrinsic region can be made long.
  - ⇒ Carrier diffusion effects minimized since all light absorbed in intrinsic region
  - ⇒ The junction capacitance is reduced compared to a pn-junction because the distance between the effective plates is increased.
  - ⇒ Carriers reach saturation velocity while traveling in intrinsic region, so even though pin depletion length  $l_p + l_d$  is longer than pn-junction depletion length, lower transit time than pn-junction where carrier velocity drops below saturation not far from metallurgical junction



### p-i-n Photodiodes

⇒ As with the pn-junction, the quantum efficiency is defined by the following equation, however the distance can now be integrated over the larger intrinsic region

$$P_{abs}(x) = P_i(1-R)(1-e^{-\alpha(\lambda)x})$$
$$= \eta(\lambda, x)P_i$$

⇒ As the depletion region length is increased,  $\eta$  increases, the junction capacitance  $C_j$  decreases, and the transit time  $\tau_{trans}$  increases. The detector design must be optimized to maximize both efficiency and bandwidth. An estimate of the bandwidth is given by

$$B_{pin} = \frac{1}{\sqrt{\left(\frac{1}{f_{RC}}\right)^2 + \left(\frac{1}{f_{trans}}\right)^2}} = \frac{1}{\sqrt{\left(2\pi R_S \varepsilon_0 \varepsilon_r \frac{A}{l_d}\right)^2 + \left(\frac{1}{0.44\nu_s}\right)^2}}$$

# Bandwidth-Efficiency Tradeoffs in p-i-n Photodiodes

- ⇒ The quantum efficiency,  $\eta$ , can be approximated assuming R=0 (high quality anti-reflection coating) and intrinsic region length  $l_d$ .  $\eta = 1 - e^{-\alpha l_d}$
- $\Rightarrow$  For small l, bandwidth is transit time limited
- ⇒ For large l, bandwidth it RC limited
- ⇒ Optimal bandwidth length where two effects are equal
- ⇒ QE keeps increasing with increased length

 $\Rightarrow If the detector area A and length <math>l_d$  are both optimized, then bandwidth and quantum efficiency can both be maximized

$$A = \frac{l_d}{2\pi R_l \varepsilon_0 \varepsilon_r} \sqrt{\frac{1}{B^2} - \left(\frac{l_d}{0.44\nu_s}\right)^2}$$
$$l_d = -\frac{1}{\alpha} \ln(1-\eta)$$



### Vertically Illuminated p-i-n Photodiodes



For a double pass vertically illuminated pin detector (see left figure), the quantum efficiency is

$$\eta = \left(1 + r e^{-\alpha l}\right) \left(1 - e^{-\alpha l}\right)$$

When the carrier transit distance is approximately equal to l, and  $\alpha l \ll 1$ , the bandwidth-efficiency for a double-pass vertically illuminated pin photodiode is approximately

# Increasing Bandwidth-Efficiency in Photodetectors

⇒ As with the pn-junction, the quantum efficiency is defined by the following equation, however the distance can now be integrated over the larger intrinsic region

### Bandwidth Efficiency Tradeoffs

### Waveguide Photodiodes

### **Distributed Photodetectors**

### **Resonant Cavity Photodetectors**

### Uni-Traveling Carrier Photodiodes

### APDs



### Staircase and Supperlattice APDs