

High-Performance InP/In_{0.53}Ga_{0.47}As/InP Double HBTs on GaAs Substrates

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Abstract—InP/In_{0.53}Ga_{0.47}As/InP double heterojunction bipolar transistors (HBTs) were grown on GaAs substrates. A 140 GHz power-gain cutoff frequency f_{max} and a 207 GHz current-gain cutoff frequency f_T were obtained, presently the highest reported values for metamorphic HBTs. The breakdown voltage V_{CEO} was 5.5 V, while the dc current gain β was 76. High-thermal-conductivity InP metamorphic buffer layers were employed in order to minimize the device-thermal resistance.

Index Terms—Heterojunction bipolar transistor, indium phosphide, metamorphic growth, molecular beam epitaxy.

I. INTRODUCTION

DOUBLE heterojunction bipolar transistors [1]–[3] (DHBTs) have applications in high-frequency communications and radar. HBTs using InP based materials (InGaAs or GaAsSb epitaxial base layers and InGaAs or InP epitaxial collector layers) currently exhibit significantly higher current-gain and power-gain cutoff frequencies than GaAs-based HBTs. However, InP substrates are expensive and fragile and are readily broken during processing. This has motivated the investigation of metamorphic growth of InP-based DHBTs on GaAs substrates [4], [5]. To date, reported metamorphic HBTs (M-HBTs) have not yet demonstrated bandwidth comparable to that of lattice-matched InP HBTs. Further, thermal performance has not yet been addressed.

Wideband DHBTs operate at high power densities. Recently reported static frequency dividers operating at 75-GHz clock frequency [6] employ InP HBTs operating at 180 kA/cm² current density and $V_{CE} \sim 1.1$ V, corresponding to a power density per unit emitter junction area A_E of $P/A_E \sim 200$ kW/cm². Current density, hence power density, must further increase as logic speed is increased [6]. W-band DHBT power amplifiers also operate at high junction power densities. As shown in the seminal work by Liu and Chau [7], [8], because of the concentrated heat flux immediately below the HBT collector, DHBT thermal resistance is critically dependent upon the thermal resistance of layers lying immediately below the DHBT collector, and thick layers of low-thermal-conductivity ($\kappa = 5$ W/k-m) InGaAs must be avoided in the subcollector. M-DHBTs face a similar difficulty, because (as measured in our laboratory) InAlAs and AlGaAsSb metamorphic buffer layers also have very low thermal conductivity (10.5 W/K-m, and 8.4 W/K-m respec-

Manuscript received January 25, 2002. This work was supported by the Office of Naval Research under Grant N00014-01-1-0065. The review of this letter was arranged by Editor D. Ritter.

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Publisher Item Identifier S 0741-3106(02)04536-6.

TABLE I
LAYER STRUCTURE OF THE MBE-GROWN InP/In_{0.53}Ga_{0.47}As/InP METAMORPHIC DHBT. ALL GRADED LAYERS ARE In_{0.53}Ga_{0.47}As/In_{0.52}Al_{0.48}As DIGITAL ALLOY GRADES, EXCEPT THE BASE, WHICH IS AN In_xGa_{1-x}As LINEAR COMPOSITIONAL GRADE. THE BASE-COLLECTOR GRADE HAS A 15 Å PERIOD

Layer	Material	Doping	Thickness (Å)
Emitter cap	In _{0.53} Ga _{0.47} As	2×10^{19} cm ⁻³ : Si	300
Grade	In _{0.53} Ga _{0.47} As to In _{0.53} Ga _{0.26} Al _{0.21} As	2×10^{19} cm ⁻³ : Si	200
N ⁺ emitter	InP	2×10^{19} cm ⁻³ : Si	700
N ⁻ emitter	InP	8×10^{17} cm ⁻³ : Si	500
Emitter-base grade	In _{0.53} Ga _{0.26} Al _{0.21} As to In _{0.455} Ga _{0.545} As	4×10^{17} cm ⁻³ : Si 8×10^{17} cm ⁻³ : Be	233 47
Graded base	In _{0.455} Ga _{0.545} As to In _{0.53} Ga _{0.47} As	4×10^{19} cm ⁻³ : Be	400
Setback	In _{0.53} Ga _{0.47} As	2×10^{16} cm ⁻³ : Si	100
Base-collector grade	In _{0.53} Ga _{0.47} As to In _{0.53} Ga _{0.26} Al _{0.21} As	2×10^{16} cm ⁻³ : Si	240
Pulse doping	InP	5.6×10^{18} cm ⁻³ : Si	30
Collector	InP	2×10^{16} cm ⁻³ : Si	1,630
Subcollector	In _{0.53} Ga _{0.47} As	1×10^{19} cm ⁻³ : Si	250
Subcollector	InP	2×10^{19} cm ⁻³ : Si	750
Buffer	InP	undoped	15,000
GaAs (100) semi-insulating substrate			

tively). These layers are typically 1.5 μ m thick and lie immediately below the DHBT subcollector.

Here we report InP-based DHBTs grown on GaAs using InP as the metamorphic buffer layer. This layer has a measured 16.1 W/K-m thermal conductivity. The current-gain and power-gain cutoff frequencies are $f_T = 207$ GHz and $f_{max} = 140$ GHz, both records for M-DHBTs. An M-DHBT with 0.4 μ m \times 7.5 μ m emitter area has a measured 2.8 K/mW thermal resistance, and exhibits a small measured 35 K junction-ambient temperature rise even when biased at $J_e = 190$ kA/cm² and $V_{ce} = 2$ V ($P/A_E = 380$ kW/cm²). At $V_{CE} = 1.5$ V, f_T and f_{max} collapse due to the Kirk effect (collector field screening) occurs at a high $J_e = 450$ kA/cm² current density.

II. GROWTH

InP/In_{0.53}Ga_{0.47}As/InP DHBTs were grown on a GaAs substrate using a Varian Gen II molecular beam epitaxy (MBE) system. Growth was initiated with a 1000 Å undoped GaAs buffer layer, followed by growth at 470 °C of a 1.5 μ m undoped InP metamorphic buffer layer. The remaining HBT layers were then grown. Key features of the layer structure (Table I) include an InP emitter, a 400 Å InGaAs base with 52 meV band-gap

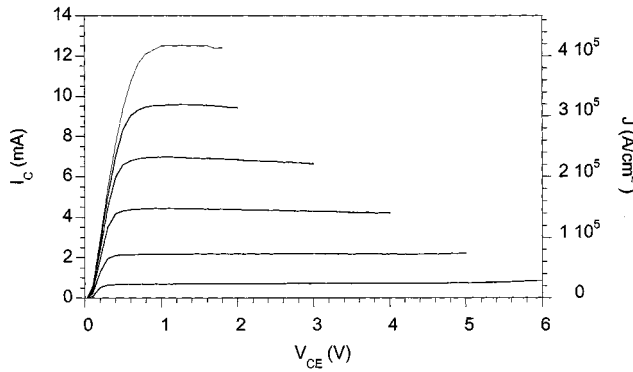


Fig. 1. Common emitter dc characteristics of $0.4 \mu\text{m} \times 7.5 \mu\text{m}$ emitter device. The base-current steps are $40 \mu\text{A}$.

grading for base transit time reduction, a 240 \AA InGaAs/InAlAs base-collector heterojunction grade, and a 1660 \AA InP collector. The heterojunction between the $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ base and the InP collector is a 240 \AA $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}/\text{In}_{0.52}\text{Al}_{0.48}\text{As}$ chirped superlattice whose composition adjacent to the collector is $\text{In}_{0.53}\text{Ga}_{0.26}\text{Al}_{0.21}\text{As}$, chosen so as to eliminate discontinuities in the conduction-band energy at the interface. A similar superlattice grade is employed in the emitter-base junction. To reduce sensitivity to diffusion of the Be base dopant, a 100 \AA undoped $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ setback layer was introduced between the base and the base-collector grade. The total collector-base depletion region thickness is therefore 2000 \AA . The metamorphically grown HBT layers are relatively rough (9.5 nm RMS , as determined by atomic force microscopy).

III. FABRICATION

HBTs were fabricated in a triple-mesa process using optical projection lithography and selective wet chemical etching. Use of narrow emitter-base and collector-base junctions reduces both the base resistance and the collector-base capacitance [9]. While the emitter contact metal is $0.7 \mu\text{m} \times 8 \mu\text{m}$, lateral undercutting during the HCl-based etch of the InP emitter forms an emitter-base junction whose dimensions are approximately $0.4 \mu\text{m} \times 7.5 \mu\text{m}$. Collector-base capacitance is reduced by employing narrow base Ohmic contacts of $0.25 \mu\text{m}$ width on either side of the emitter stripe, producing a small $1.2 \mu\text{m} \times 11 \mu\text{m}$ base-collector junction area. Polyimide is used both for passivation and for mesa planarization prior to interconnect deposition.

IV. RESULTS

Fig. 1 shows the common emitter characteristics, measured from 0 – 400 kA/cm^2 current density. The measured dc current gain is approximately 76, the common-emitter open-circuit breakdown voltage at low current densities BV_{CEO} is greater than 6 V, while $V_{\text{CE,SAT}} < 0.8 \text{ V}$ at 400 kA/cm^2 current density. The measured sheet resistance of the metamorphic InP buffer layer is approximately $150 \text{ M}\Omega$ per square. The base sheet resistance is 1020Ω per square. This sheet resistance was measured after the emitter mesa etch, which removes approximately the upper 100 \AA of the base layer. The corresponding hole mobility is $51 \text{ cm}^2 \text{ cm/V-s}$, 10% lower than we observed for similar lattice-matched DHBT growths.

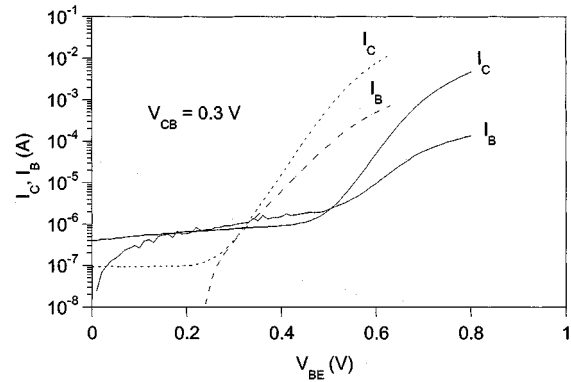


Fig. 2. Metamorphic HBT Gummel characteristics of (solid lines) an HBT with a $0.4 \mu\text{m} \times 7.5 \mu\text{m}$ emitter-base junction and a $1.2 \mu\text{m} \times 11 \mu\text{m}$ base-collector junction, and an HBT (dotted lines) with a $60 \mu\text{m} \times 60 \mu\text{m}$ emitter-base junction and a $100 \mu\text{m} \times 130 \mu\text{m}$ base-collector junction.

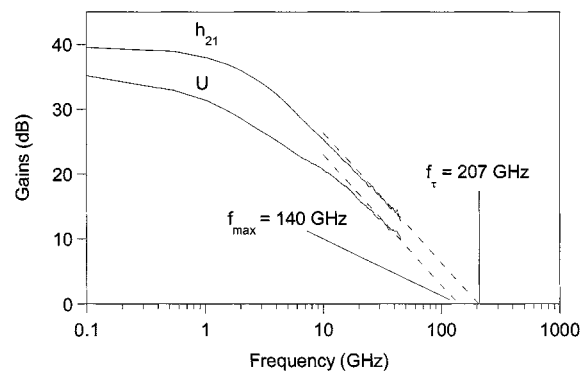


Fig. 3. Measured short-circuit current gain h_{21} and Mason's unilateral power gain U versus frequency for an HBT with a $0.4 \mu\text{m} \times 7.5 \mu\text{m}$ emitter-base junction and a $1.2 \mu\text{m} \times 11 \mu\text{m}$ base-collector junction. $I_C = 12.0 \text{ mA}$ and $V_{\text{CE}} = 1.5 \text{ V}$.

Fig. 2 shows the HBT Gummel [$\log(I_C, I_B)$ versus V_{CE}] characteristics, indicating a collector current ideality factor of 1.1 and a base current ideality factor of 2.2. The characteristics are measured with nonzero (0.3 V) reverse bias applied to the collector-base junction, so that base-collector junction leakage, if present, will be observed. For comparison, Fig. 2 also shows the Gummel characteristics of a large-area HBT ($60 \mu\text{m} \times 60 \mu\text{m}$ emitter-base and $100 \mu\text{m} \times 130 \mu\text{m}$ base-collector junctions) fabricated from the same epitaxial material. Despite the large collector-base junction area, the Gummel characteristics indicate that I_{cbo} of the large-area HBT is below $0.1 \mu\text{A}$, indicating a collector-base leakage current per unit junction area below 10^{-3} A/cm^2 . The small-area HBTs are fabricated on a mask set designed for microwave on-wafer probing calibrated using the line-reflect-line method, and have long ($230 \mu\text{m}$) 50Ω coplanar waveguide (CPW) transmission lines between the probe pads and the transistor terminals. The leakage currents observed at low V_{be} in Fig. 2 for the small-area HBT are dominated by leakage currents associated with conduction through the buffer layer between these long interconnects. HBTs are minority-carrier devices, and defect-induced leakage currents resulting from metamorphic growth are potential concerns. Nevertheless, high dc current gains and low $< 10^{-3} \text{ A/cm}^2$ base-collector leakage are observed.

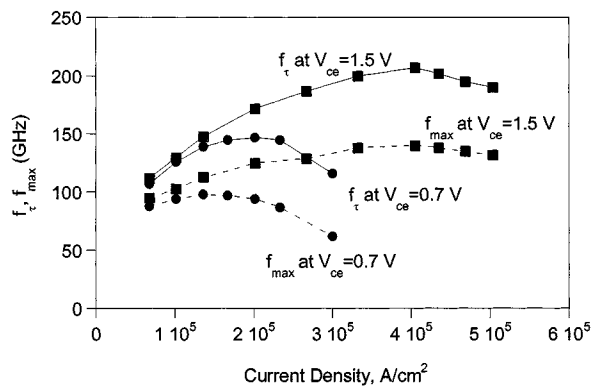


Fig. 4. Measured current-gain cutoff frequency f_{τ} and power-gain cutoff frequency f_{\max} versus current density at $V_{CE} = 0.7$ V and $V_{CE} = 1.5$ V.

Fig. 3 shows the current gain (h_{21}) and unilateral power gain (U) of the small-area HBT, computed from the measured 0.045–45 GHz S -parameters. A 207 GHz f_{τ} and a 140 GHz f_{\max} were measured at $I_C = 12.0$ mA (4×10^5 A/cm²) and $V_{CE} = 1.5$ V, as determined by a -20 dB/decade extrapolation. These are the highest values reported for metamorphic HBTs. Fig. 4 shows the variation of f_{τ} with collector current density, as measured at $V_{CE} = 0.7$ V and $V_{CE} = 1.5$ V. The observed decrease in f_{τ} at very high current densities (2.5×10^5 A/cm² at $V_{CE} = 0.7$ V, 4.5×10^5 A/cm² at $V_{CE} = 1.5$ V) is due to the Kirk effect. HBT junction temperature and thermal resistance were measured by the method of [7]. The $0.4 \mu\text{m} \times 7.5 \mu\text{m}$ emitter HBTs have a 2.8 K/mW thermal resistance and exhibit a small measured 35 K junction-ambient

temperature rise when biased at $J_e = 190$ kA/cm² and $V_{ce} = 2$ V ($P/A_E = 380$ kW/cm²).

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